String sketches of Dynamical supersymmetry breaking

M. Bíanchí, F.Fucíto, J.F.M, hep-th/0904.2156

F. Fucito, A. Lionetto, J.F.M, R. Richter, work in progress



• D-brane worlds:

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MSSM particles: Open strings
 (orbifolds, intersecting or magnetized branes)

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Interactions: String diagrams (gauge and Yukawa couplings, instantons) **D**-branes



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 $R^{1,3} \times CY$

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Non-perturbative superpotentials: + D- and F- term interactions (fluxes, gauge/gravity mediation from hidden sectors)

MSSM (1 or 3 generations)+hidden SQCD (intersecting branes)

SU(5)+2 generations (orbifold quiver)

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Consistent models: with one or three MSSM generations

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$$\Longrightarrow \xi \neq 0$$
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- Conclusions:

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\Box_a	$\frac{1}{2}\left(\pi_a \circ \pi'_a + \pi_a \circ \pi_{\rm O6}\right)$
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✤ Tadpole conditions:

$$\sum_{a} N_a \left(\pi_a + \pi'_a \right) - 4\pi_{O6} = 0$$

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Anomaly cancelation + ...

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$$\sum_{a} N_a (\pi_a + \pi'_a) - 4\pi_{O6} = 0$$

$$\Rightarrow \Delta n(\Box_a) + (N_a + 4)\Delta n(\Box_a) + (N_a - 4)\Delta n(\Box_a) = 0$$
Anomaly cancellation + ...
$$\Rightarrow \text{Masslessness of } U(1)_Y : \sum_{x} q_x N_x(\pi_x - \pi'_x) = 0$$

$$U(1)_Y = \sum_{x} q_x U(1)_x$$

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$$\implies q_a N_a \left(\Delta n(\Box_a) + \Delta n(\Box_a)\right) + \sum_{x \neq a} q_x N_x \Delta n(\Box_a, \Box_x) = 0$$

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Field theory Models

Field theory Models

• Gauge group: $G = U(3)_a \times SU(2)_b \times U(1)_c \times U(1)_d \times U(1)_e \left(U(3)_f \right)_f$

MSSM+extra U(1)'s

5

Hidden SQCD

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• Hypercharge: $U(1)_Y = \frac{1}{6}U(1)_a + \frac{1}{2}\left[U(1)_c + U(1)_d + U(1)_e + U(1)_f\right]$

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♦ Hidden sector: U(3) SQCD+2 quark-antiquark pairs $W_{np} = \frac{\Lambda^7}{\det(M_{ij})}$ Mesons: $M_{ij} = Q_i \cdot \tilde{Q}_j$

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Field Content:
Quarks and Leptons: Higgs and Mesons:

$$\vec{\Phi} = \{Q_L, L, u^c, d^c, e^c, \nu^c\}$$
$$\vec{X} = \{H_u, H_d, M_i\}$$

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Quarks and Leptons: $\vec{\Phi} = \{Q_L, L, u^c, d^c, e^c, \nu^c\}$ Higgs and Mesons: $\vec{X} = \{H_u, H_d, M_i\}$ ♦ N=1 Lagrangian: $\mathcal{L}_{D} = \left(k_{ij}(\vec{X}, \vec{X}^{\dagger}) \Phi_{i}^{\dagger} e^{V} \Phi_{j} + \vec{X}^{\dagger} e^{V} \vec{X} + \xi_{a} V_{a} \right) \Big|_{\theta^{2} \bar{\theta}^{2}}$ non-renormalizable $\mathcal{L}_F = \left(\tau_a(\vec{X}) \operatorname{Tr} W^{(a)} W^{(a)} + W(\vec{X}, \vec{\Phi}) \right) \Big|_{a^2}$

Jose F. Morales (INFN, Tor Vergata, Roma)

Gauge group: $G = U(3)_a \times SU(2)_b \times U(1)_c \times U(1)_d \times U(1)_e$ (U(3)) MSSM+extra U(1)'s Hidden SQCD Hypercharge: $U(1)_Y = \frac{1}{6}U(1)_a + \frac{1}{2}[U(1)_c + U(1)_d + U(1)_e + U(1)_f]$ ♦ Hidden sector: U(3) SQCD+2 quark-antiquark pairs $W_{np} = \frac{\Lambda'}{\det(M_{ij})}$ Mesons: $M_{ij} = Q_i \cdot \tilde{Q}_j$ ♦ Field Content:
Quarks and Leptons: $\vec{\Phi} = \{Q_L, L, u^c, d^c, e^c, \nu^c\}$ Higgs and Mesons: $\vec{X} = \{H_u, H_d, M_i\}$ ♦ N=1 Lagrangian: $\mathcal{L}_{D} = \left(k_{ij}(\vec{X}, \vec{X}^{\dagger}) \Phi_{i}^{\dagger} e^{V} \Phi_{j} + \vec{X}^{\dagger} e^{V} \vec{X} + \xi_{a} V_{a} \right) \left| \underbrace{\leftarrow}_{\theta^{2} \bar{\theta}^{2}} \text{Charged Higgs and Mesons} \right|_{\theta^{2} \bar{\theta}^{2}}$ $\mathcal{L}_F = \left(\tau_a(\vec{X}) \operatorname{Tr} W^{(a)} W^{(a)} + W(\vec{X}, \vec{\Phi}) \right) \Big|_{a^2}$ Yukawa int for Higgs and Mesons 5 Jose F. Morales (INFN, Tor Vergata, Roma) Dynamical susy breaking Perugia 2010

• The prepotential: $W(\vec{X},\vec{\Phi}) = W_0(\vec{X}) + W_2(\vec{X},\vec{\Phi})$

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$$W_0 = \mu H_u H_d M_1 + m M_1 M_2 + \frac{\Lambda^3}{M_1 M_2}$$

	H_u	H_d	M_1	M_2
SU(2)	2	2	1	1
$U(1)_d$	0	-1	1	-1
$U(1)_e$	1	0	-1	1

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The prepotential: $W(\vec{X}, \vec{\Phi}) = W_0(\vec{X}) + W_2(\vec{X}, \vec{\Phi})$ **MSSM** Yukawa interactions susy and gauge breaking we look for vacua with : $\langle \vec{\Phi} \rangle = 0$ $\langle \vec{X} \rangle = \vec{x} + \theta^2 \vec{F}_x$ $W_0 = \mu H_u H_d M_1 + m M_1 M_2 + \frac{\Lambda^5}{M_1 M_2}$ H_u H_d M_1 M_2 SU(2)2 2 1 1 $U(1)_{d}$ 0 _1 1 -1non-perturbative Higgs-Meson interactions $U(1)_e$ 1 0 -11 The potential : $V = |\vec{F}|^2 + \frac{1}{2}\vec{D}^2$



• The susy vacuum: $\xi_d + \xi_e = 0$

$$h_u = h_d = 0 \qquad M_1 = M_2 = \left(\frac{\Lambda^5}{m}\right)^{\frac{1}{4}}$$

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• Non supersymmetric vacua : $\xi_d + \xi_e \neq 0$

Gauge unbroken solution:

$$g_b = g_d = g_e = 0.5 \qquad m = \mu = \Lambda = \xi_e = \xi_d = 1$$

$$h_u = h_d = 0 \qquad x_1 = x_2 = 1$$

$$V = \frac{g_d^2 g_e^2 (\xi_d + \xi_e)^2}{2(g_d^2 + g_e^2)}$$

$$V = 0.25$$

♦ Gauge broken solution:

 $g_b = g_d = g_e = \Lambda = 0.5$ $m = \mu = \xi_e = \xi_d = 1$ $h_u = 0.01$ $h_d = 0.5$ $x_1 = 0.37$ $x_2 = 0.46$ V = 0.23



One generation model:

d **Q** ** u $\widetilde{\mathsf{Q}}_1$ Q_1 \leftarrow H_u H_d e ν $\bigstar M_1 = \tilde{Q}_1 Q_2$ $\rightarrow M_2 = \tilde{Q}_2 Q_1$











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- ★ Baryons and Mesons: $X^{I} = C^{Iu}E_{u} \quad Y^{i}_{I} = \frac{1}{2}\epsilon_{IJK}A^{i}_{u_{1}u_{2}}C^{Ju_{1}}C^{Ku_{2}}$ $\tilde{Y}^{ij} = \frac{1}{4}\epsilon^{u_{1}...u_{5}}A^{i}_{u_{1}u_{2}}A^{j}_{u_{3}u_{4}}E_{u_{5}}$ $Z^{iI} = \frac{1}{12}\epsilon^{u_{1}...u_{5}}A^{i}_{u_{1}u_{2}}A^{j}_{u_{3}u_{4}}A_{u_{5}v,j}C^{Iv}$



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Quantum deformed constraints:

$$Y_{I}^{i} Z_{i}^{I} = 0$$

 $\epsilon_{IJK} X^{I} Z_{i}^{J} Z^{iK} + Y_{iI} \tilde{Y}^{ij} Z_{j}^{I} = \Lambda^{10}$



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Susy

Pe



Conclusions

Jose F. Morales (INFN, Tor Vergata, Roma) Dynamical susy breaking





Brane intersections, orbifold singularities: F-, D- type interactions+instantons





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