#### **BEAUTY 2005**

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#### **Cecilia Tarantino**

Università Roma Tre and INFN Sezione di Roma III



1. Mixing Δm<sub>d</sub>, Δm<sub>s</sub>

- Phenomenological importance (UTA)
- Theoretical inputs







• $\eta_B S_0(x_t)$  : short-distance physics, from Perturbation Theory (2%) [at NLO, Buras et al.: Nucl.Phys.B347 (1990)]

•  $\hat{B}_{B_s} f_{B_s}^2, \xi^2$  : non-perturbative QCD effects, from Lattice QCD

: lattice uncertainties cancel



### **Lattice QCD Actions:**

•Wilson/O(a)-improved Wilson: problem at light masses

(exceptional quenched configurations)

•**Twisted mass:** no exceptional configurations, easier improving, but more expensive •**Staggered:** light masses, but 4 tastes (fourth root trick)

•Domain-wall/Overlap: light masses, but the most expensive

Accurate "quenched" studies + recent Nf=2, Nf=2+1 calculations





### **B-B** Mixing on the Lattice



JLQCD(2003): N<sub>f</sub>=2 & NRQCD



Lattice data are consistent with a constant.

## **B** Mixing from Lattice QCD

	ICHEP 2002 (Lellouch)	BEAUTY 2005 (ICHEP 2004, Hashimoto)
$f_{B_d}$ (MeV)	203(27)(+0-20)	189(27)
$f_{B_s}$ (MeV)	238(31)	230(30)
$f_{B_s}\sqrt{\hat{B}_{B_s}}$ (MeV)	276(38)	262(35)
$f_{\scriptscriptstyle B_s}$ / $f_{\scriptscriptstyle B_d}$	1.18(4)(+12-0)	1.22(+5-6)
ξ	1.18(4)(+12-0)	1.23(6)

Now averages include rough ``estimates'' of chiral logs (m<sub>q</sub>/m<sub>s</sub>>0.2) and unquenched effects (N<sub>f</sub>=2+1)

## **B Mixing from QCD Sum Rules**

- 1. Equating phenomenological and theoretical spectral functions;
- 2. Determination of theoretical spectral functions by calculating two or three-point correlators in perturbative QCD, including corrections from the OPE.

 $O(\alpha_s^2)$  recently calculated  $O(\alpha_s)$  completed but non-factorisable contributions important

$$f_B = 210(19) \text{ MeV},$$
  
 $f_{B_s} = 244(21) \text{ MeV}$  (Jamin et al '02)  
 $\hat{B}_B = 1.60(3)$  (Körner et al '03)

 $f_{\mathsf{B}}$ 

**Problem: many parameters, loosely constrained!** 



$$\Gamma_{21}^q \quad (\Delta B = 2)$$

•Width Differences  $\Delta\Gamma_{B_d}$ ,  $\Delta\Gamma_{B_s}$ 

•CP-Violation Parameters  $|q/p|_{B_{d,s}}$ 

$$\Gamma_{11}^q \quad (\Delta B = 0)$$

•Lifetime Ratios  $\tau(B^+)/\tau(B_d), \tau(B_s)/\tau(B_d), \tau(\Lambda_b)/\tau(B_d)$ 



Wilson Coefficients from P.T.

•**Lifetime Ratios**  $\Gamma = \frac{G_F^2 |V_{cb}|^2 m_b^5}{192 \pi^3 (2M_B)} \left[ c^{(3)} \langle \overline{b} b \rangle + c^{(5)} \frac{g_s}{m_b^2} \langle \overline{b} \sigma_{\mu\nu} G^{\mu\nu} b \rangle + \frac{96 \pi^2}{m_b^3} \sum_k \left( c_k^{(6)} \langle O_k^{(6)} \rangle + \frac{c_k^{(7)}}{m_b} \langle O_k^{(7)} \rangle \right) \right]$  *O*(1) (1996) [M. Neubert and C.T. Sachrajda] *O*(\alpha\_s) (2002)

•[E. Franco, V. Lubicz, F. Mescia and C.T.] •[M. Beneke, G. Buchalla, C. Greub, A. Lenz and U. Nierste]

 $O(1/m_b)$  (2004) •[F. Gabbiani, A. I. Onishchenko and A. A. Petrov]

•Width Differences and 
$$\Gamma_{21}^{s} = -\frac{G_{F}^{2}m_{b}^{2}}{12\pi(2M_{B_{s}})}(V_{cb}^{*}V_{cs})^{2}[G\langle Q \rangle + G_{S}\langle Q_{S} \rangle + \delta_{1/m_{b}}]$$

•[M. Ciuchini, E. Franco, V. Lubicz, F. Mescia and C.T.]

•[M. Beneke, G. Buchalla, A. Lenz and U. Nierste]

 $O(1/m_b)$  (1996) •[M. Beneke, G. Buchalla and I. Dunietz]

# **Matrix Elements:** $\Delta B = 2$ **Operators**

**Leading contribution**  $O(1/m_b^3)$ 

$$\mathcal{O}_1^q = (\bar{b}q)_{V-A} \, (\bar{b}q)_{V-A} \leftrightarrow B_1^q, \quad \mathcal{O}_2^q = (\bar{b}q)_{S-P} \, (\bar{b}q)_{S-P} \leftrightarrow B_2^q.$$

 $B_q^1, B_q^2$ From the *lattice* (with different methods)<br/>or QCD -sum rules [J.G. Korner et al., 2003]Subleading contribution $O(1/m_b^4)$ 

**4 operators**  $(R_1^q, R_2^q, R_3^q, R_4^q)$ :

- $R_1^q, R_4^q$ : related, through Fierz and eq. of motion, to operators computed on the *lattice*
- $R_2^q, R_3^q$  : from the VSA

 $B_1^s = 0.87(2)(5), \ B_2^s = 0.84(2)(4)$ 

[APE (D. Becirevic et al.), 2001]

•<u>QCD +HQET</u>

 $B_1^s = 0.91(3)^{+0}_{-6}, \ B_2^s = 0.86(2)^{+2}_{-3}$ 

[APE (D. Becirevic et al.), 2000]

•**QCD**  $(m_c \leq m_Q < m_b, m_Q \rightarrow m_b)$ 

 $B_1^s = 0.85(2)(6), \ B_2^s = 0.84(6)(8) \quad \text{[JLQCD (S. Aoki et al.), 2001-2003]}$ 

•*unquenched* NRQCD  $n_f = 2$ 

 $B_1^s = 0.85(3)(11), \ B_2^s = 0.82(2)(11)$  [Hi-KEK (S. Hashimoto et al.), 2000]

•**NRQCD**  $\mathcal{O}(1/m_b)$ 

 $B_1^s = 0.83(5)(6), \ B_2^s = 0.81(2)(10)$  [V. Gimenez and J. Reyes, 2000]

•HQET  $(m_b \to \infty)$  on the lattice



#### SPQcdR, 2001



 Combine the static HQET results for *B*-parameters with the relativistic lattice QCD ones
⇒ extrapolation → "interpolation"

 Perturbative matching of the anomalous dimensions of 4-f QCD and HQET operators made @ NLO in perturbation theory!

• So far, the approach has been only applied in the quenched case

### **<u>Matrix Elements</u>**: $\Delta B = 0$ **Operators**

**Leading spectator effect contribution**  $O(1/m_b^3)$ 

$$\begin{split} \mathcal{O}_1^q &= (\bar{b}q)_{V-A} \, (\bar{q}b)_{V-A} \leftrightarrow \mathbf{B}_1^q, \qquad \mathcal{O}_2^q &= (\bar{b}q)_{S-P} \, (\bar{q}b)_{S+P} \leftrightarrow \mathbf{B}_2^q, \\ \mathcal{O}_3^q &= (\bar{b}T^a q)_{V-A} \, (\bar{q}T^a b)_{V-A} \leftrightarrow \mathbf{\epsilon}_1^q, \quad \mathcal{O}_4^q &= (\bar{b}T^a q)_{S-P} \, (\bar{q}T^a b)_{S+P} \leftrightarrow \mathbf{\epsilon}_2^q, \end{split}$$

$$O_P = (\bar{b}T^a b)_V \sum_{q=u,d,s,c} (\bar{q}T^a q)_V \qquad (\text{not computed})$$

 $\left[ (\overline{q}q)_{V-A} = \overline{q}\gamma_L^{\mu}q, \ (\overline{q}q)_{S\pm P} = \overline{q}(1\pm\gamma_5)q, \ (\overline{q}q)_V = \overline{q}\gamma^{\mu}q \right]$ 



$$\begin{array}{c|c} B_{d} - B_{s} - B^{+} & \begin{tined} \hline B-parameters \\ \hline B_{d} - B_{s} - B^{+} & \begin{tined} \hline B-parameters \\ \hline B-$$





**Theoretical predictions at the NLO + contribution of**  $O(1/m_b^4)$  :

$$\frac{\Delta\Gamma_d}{\Gamma_d} = (2.42 \pm 0.59)10^{-3} \qquad \frac{\Delta\Gamma_s}{\Gamma_s} = (7.4 \pm 2.4)10^{-2}$$

[M. Ciuchini, E. Franco, V. Lubicz, F. Mescia and C. T., 2003]

#### **Experimental measurements:**

$$\frac{\Delta\Gamma_d}{\Gamma_d} = 0.008 \pm 0.037 \pm 0.018$$
 (BaBar collaboration, 2003)

$$\frac{\Delta\Gamma_s}{\Gamma_s} = 0.07^{+0.09}_{-0.07}$$
 (HFAG, 2004)



$$c_1 << c_2$$

$$\frac{c_2^{(1)}}{c_2^{(0)}} = -27\%$$
 (large cancellations at the NLO)

 $\frac{\delta_{1/m_b}}{c_2 \langle O_2^s \rangle} = -55\% \quad \text{(large cancellations at } O(1/m_b^4))$ 

**CP-Violation Parameters**  $|\binom{q}{p}_{q}| - 1 = \frac{1}{2}\mathcal{I}m(\frac{\Gamma_{21}^{q}}{M_{21}^{q}})$ 

 $(|(q/p)_q| - 1)/\Delta\Gamma_q = \mathcal{O}(m_c^2/m_b^2), (|(q/p)_s| - 1)/(|(q/p)_d| - 1) = \mathcal{O}(\lambda^2)$ 



**Theoretical predictions at the NLO:** 

 $|(q/p)_d| - 1 = (2.96 \pm 0.67) 10^{-4}$ ,  $|(q/p)_s| - 1 = -(1.28 \pm 0.27) 10^{-5}$ 

[M. Ciuchini, E. Franco, V. Lubicz, F. Mescia and C. T., 2003]

**Experimental measurement:** 

 $|(q/p)_d| - 1 = 0.0013 \pm 0.0034$ 

[HFAG, 2005]

•Importance of more accurate measurements > UTA

**Lifetime Ratios** 



NLO









	$\frac{\tau(B^+)}{\tau(B_+)}$	$\frac{\tau(B_s)}{\tau(B_s)}$	$\frac{\tau(\Lambda_b)}{\tau(B_b)}$
LO	1.01(3)	<b>1.00</b> (1)	0.93(4)
NLO	1.06(3)	1.00(1)	0.90(5)
NLO+	1.06(2)	1.00(1)	0.88(5)
$O(1/m_b^4)$			

<u>Theoretical predictions at the NLO + contribution of  $O(1/m_h^4)$ :</u>

$$\frac{\tau(B^+)}{\tau(B_d)} = 1.06 \pm 0.02, \ \frac{\tau(B_s)}{\tau(B_d)} = 1.00 \pm 0.01, \ \frac{\tau(\Lambda_b)}{\tau(B_d)} = 0.88 \pm 0.05$$

[E. Franco, V. Lubicz, F. Mescia and C. T., 2002-2003]

 $\tau(\Lambda_h)/\tau(B_d)$  at 10

**Experimental measurements:** 

 $\frac{\tau(B^+)}{\tau(B_d)} = 1.081 \pm 0.015, \ \frac{\tau(B_s)}{\tau(B_d)} = 0.939 \pm 0.044, \ \frac{\tau(\Lambda_b)}{\tau(B_d)} = 0.803 \pm 0.047$ 

[LEP+CDF+B-factories average, Heavy Flavor Averaging Group (HFAG), 2004]

•Good agreement at the NLO and  $O(1/m_b^4)$ 

## Conclusions

**MIXING:** Phenomenological importance of  $\Delta m_d$ ,  $\Delta m_s$  in the UTA Theoretical inputs at 5-15% from Lattice QCD



