Recent Results from the Graal Beam

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for the Graal collaboration

Sixth International Conference on Perspectives in Hadronic Physics

ICTP - Trieste, May 12-16, 2008

Monday, May 26, 2008

Graal Apparatus

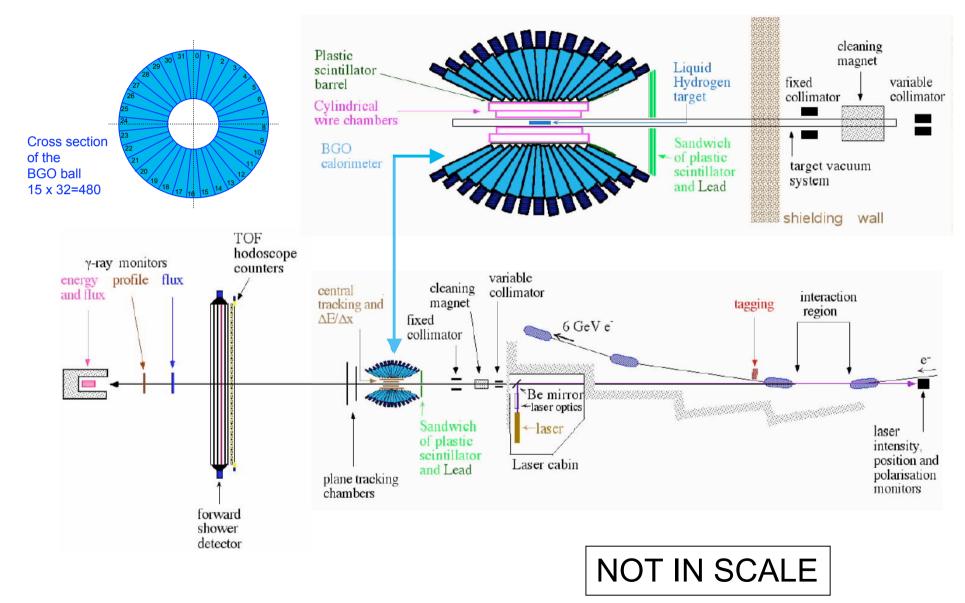
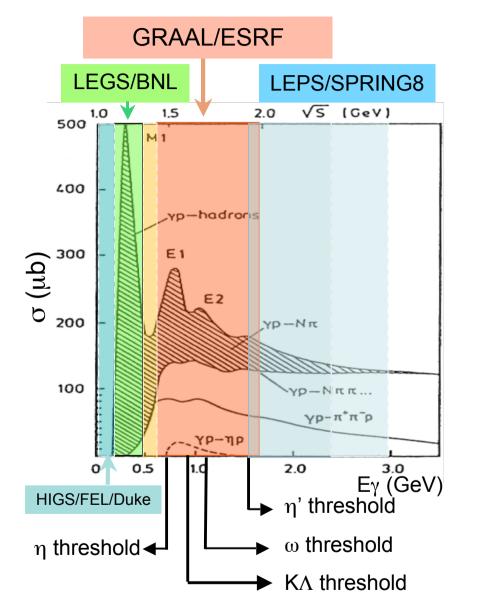


Table. Ladon beams world	dwide
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Project name	Ladon	Taladon	ROKK-1M	LEGS	LEGS-2	Graal	LEPS	HIGS
Location	Frascati, Italy		Novosibirsk, Russia	Brookhaven, US		Grenoble, France	Harima, Japan	Durham, UK
Storage ring	Adone	Adone	VEPP-4M	NSLS	NSLS	ESRF	SPring-8	TUNL-FEL
Energy defining method	Collimation	Internal tagging	Tagging	External tagging	External tagging	Internal tagging	Internal tagging	Collimation
Electron energy (GeV)	1.5	1.5	1.4–5.3	2.5	2.8	6.04	8	1.0
Laser photon energy (eV)	2.45	2.45	1.17-3.51	3.53	4.71	3.53	3.53	8.2
Gamma-ray energy (MeV)	5-80	35-80	100-1200	180–320	285-470	550-1470	1500-2400	5-225
Energy resolution (%)	1.4–10	5	_	1.6	1.1	1.1	1.25	1
Energy spread (FWHM, MeV)	0.07-8	4–2	_	5	5	16	30	_
Electron current	0.1	0.1	0.1	0.2	0.2	0.2	0.1	100
Gamma intensity s ⁻¹	105	5×10^5	2×10^{6}	4×10^{6}	2×10^{6}	2×10^{6}	2×10^{6}	10 ⁶ -10 ⁸
First year of operation	1978	1989	1993	1987	1999	1996	1999	1996

Ladon Beams in the World



•Graal:

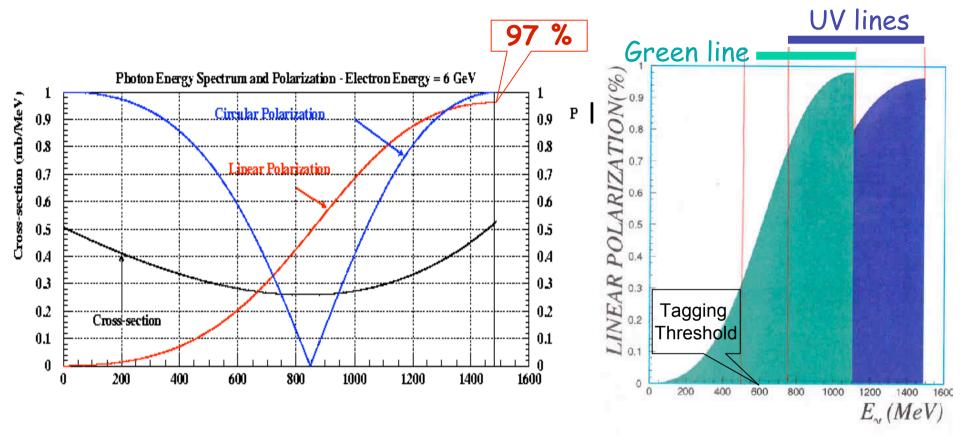
• E_{γ} = .6-1.5 GeV / W=1.4-1.9 GeV

 Region of the second and third baryon resonances

• η , K, ω , η' thresholds

 Complementary of HIGS, LEGS, Graal and LEPS

Polarization of the Graal beam



At maximum gamma-ray energy the polarization is very close to that of the laser. Changing the laser line changes the polarization of the gamma-ray beam at a given energy.

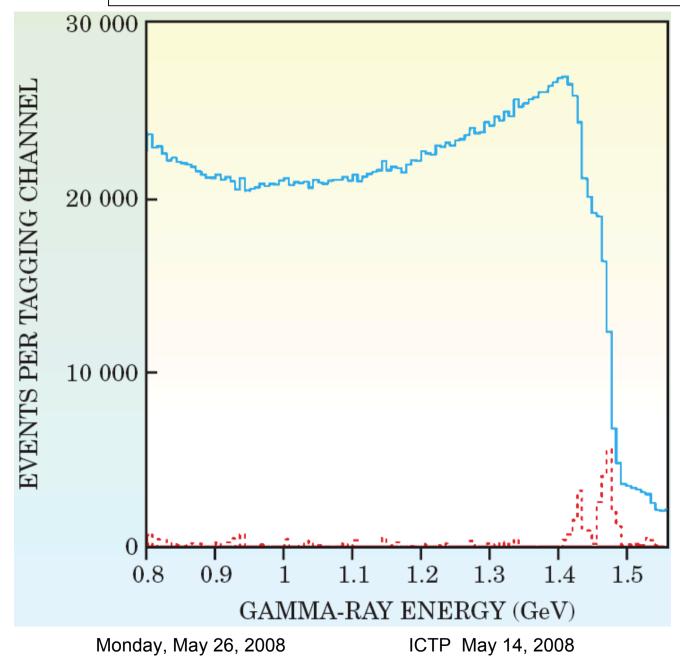
Graal Experimental Program

 $\vec{\gamma} + p \rightarrow \pi^0 + p$ $\vec{\gamma} + p \rightarrow \eta + p$ $\vec{\gamma} + p \rightarrow \pi^+ + n$ $\vec{\gamma} + \mathbf{p} \rightarrow \pi^0 + \pi^0 + \mathbf{p}$ Finished $\vec{\gamma} + \mathbf{p} \rightarrow \pi^0 + \eta + \mathbf{p}$ $\vec{\gamma} + \mathbf{p} \rightarrow k^+ + \Lambda$ $\vec{\gamma} + p \rightarrow k^{+} + \Sigma^{0}$ $\vec{\gamma} + n \rightarrow \eta + n$

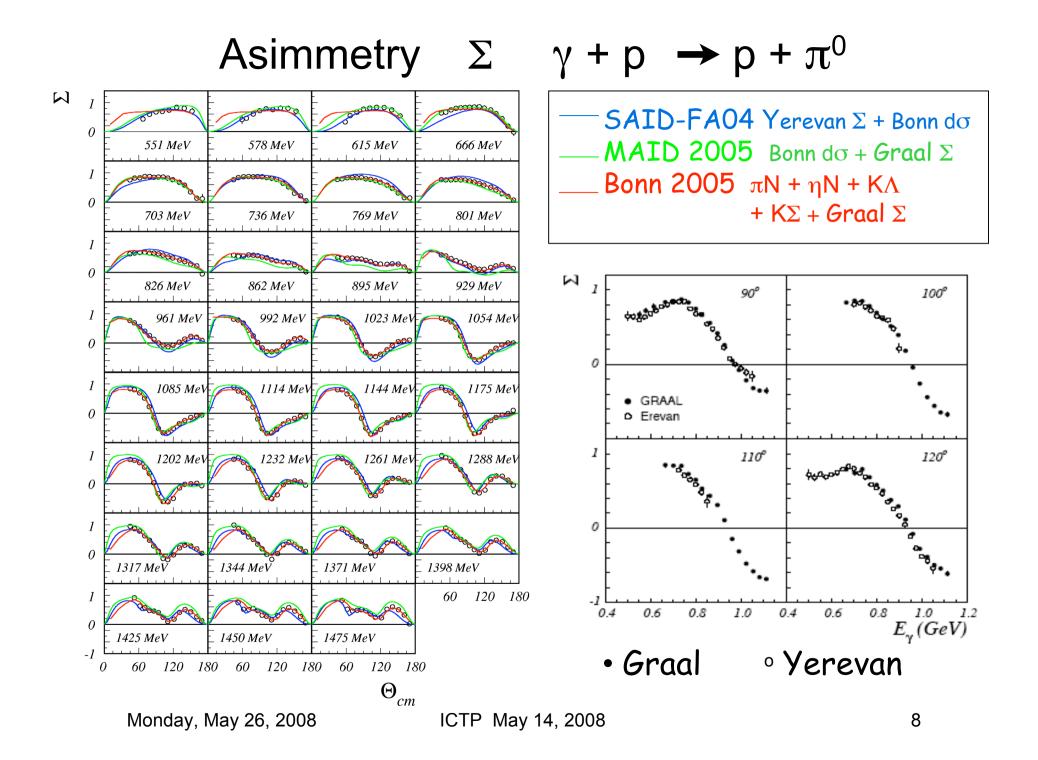
Work in progress

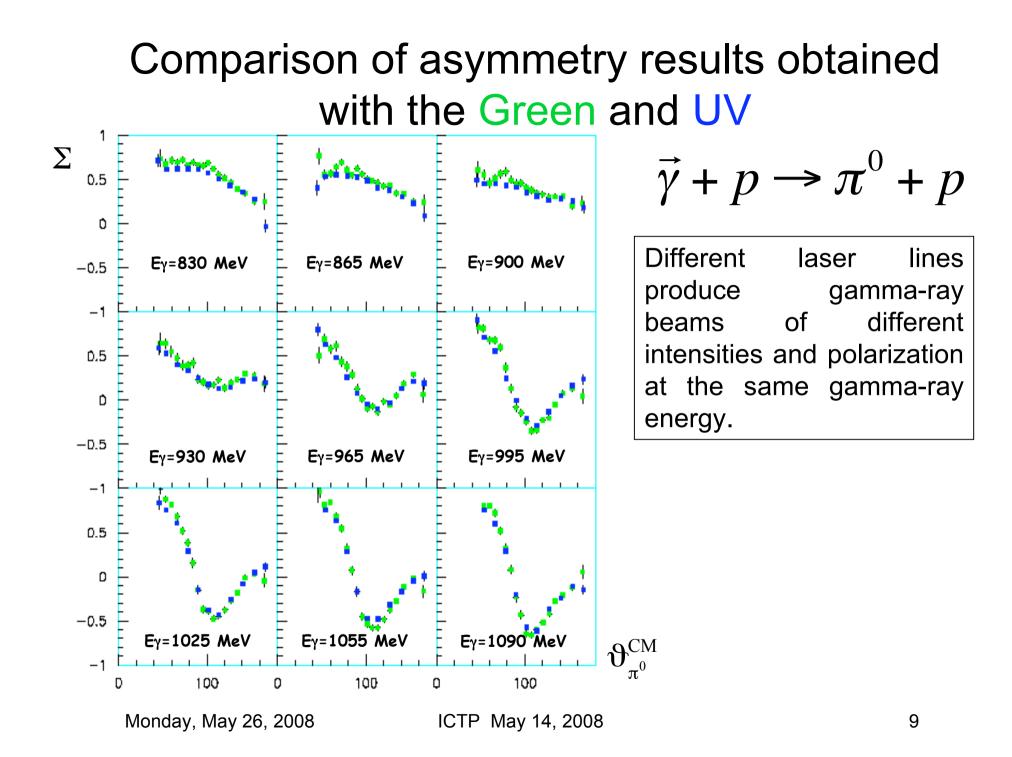
$$\vec{\gamma} + n \rightarrow \pi^{0} + n$$
$$\vec{\gamma} + p \rightarrow \omega + p$$
$$\vec{\gamma} + d \rightarrow p + n$$

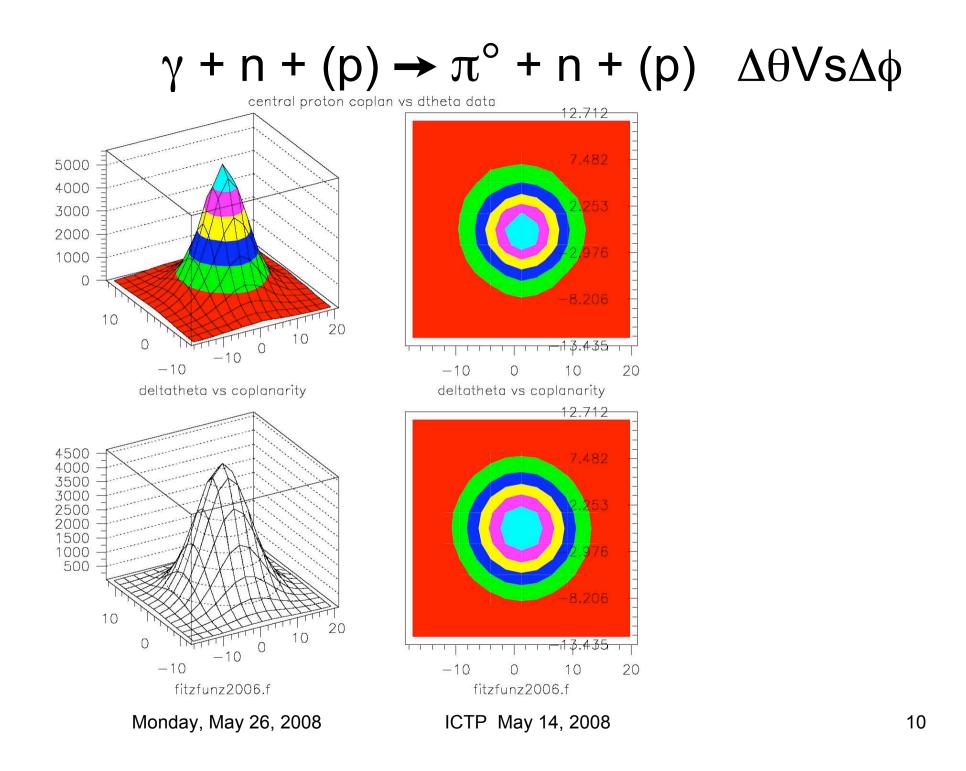
Gamma-ray Energy Spectrum of Ladon Beams

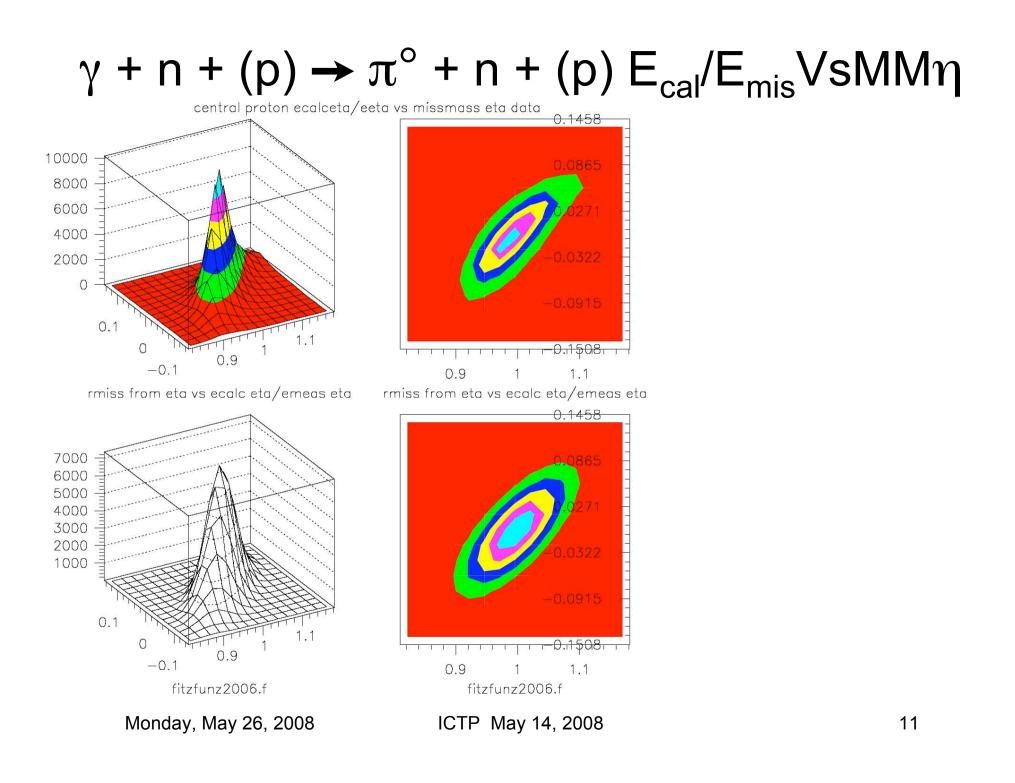


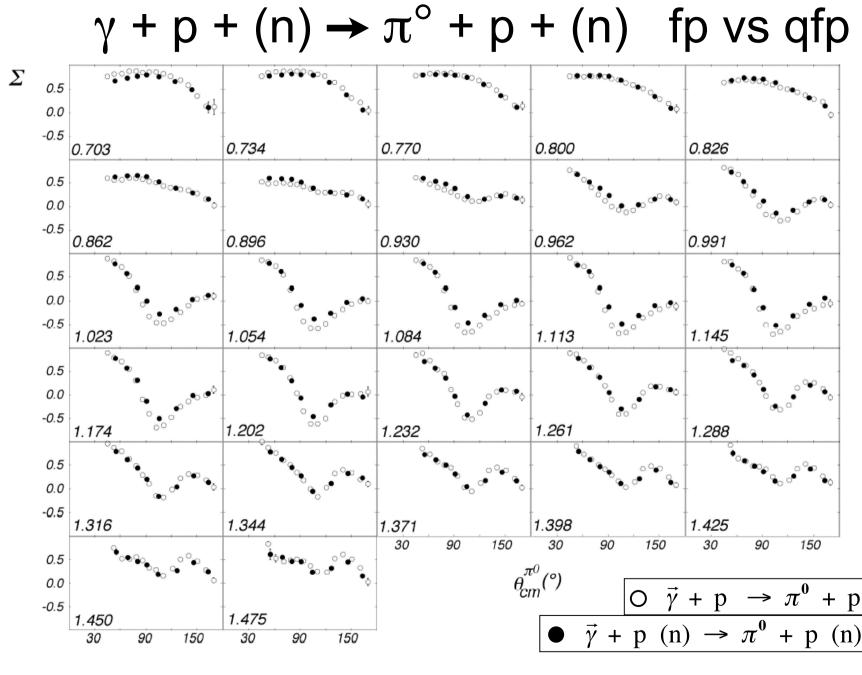
Graal beam with 3 UV laser lines

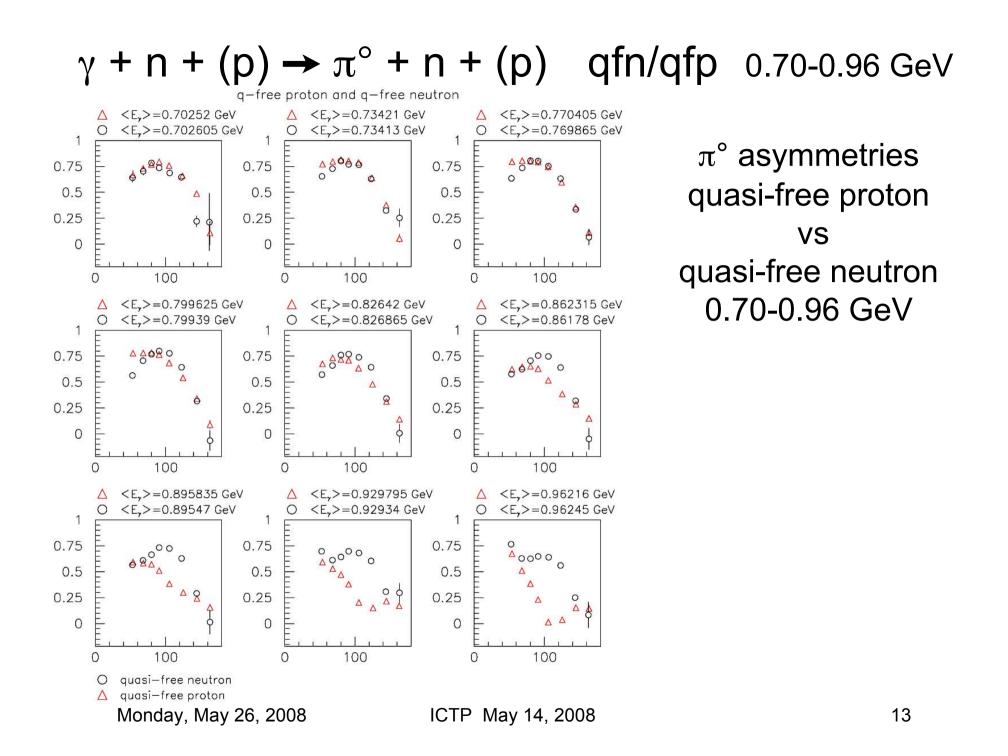


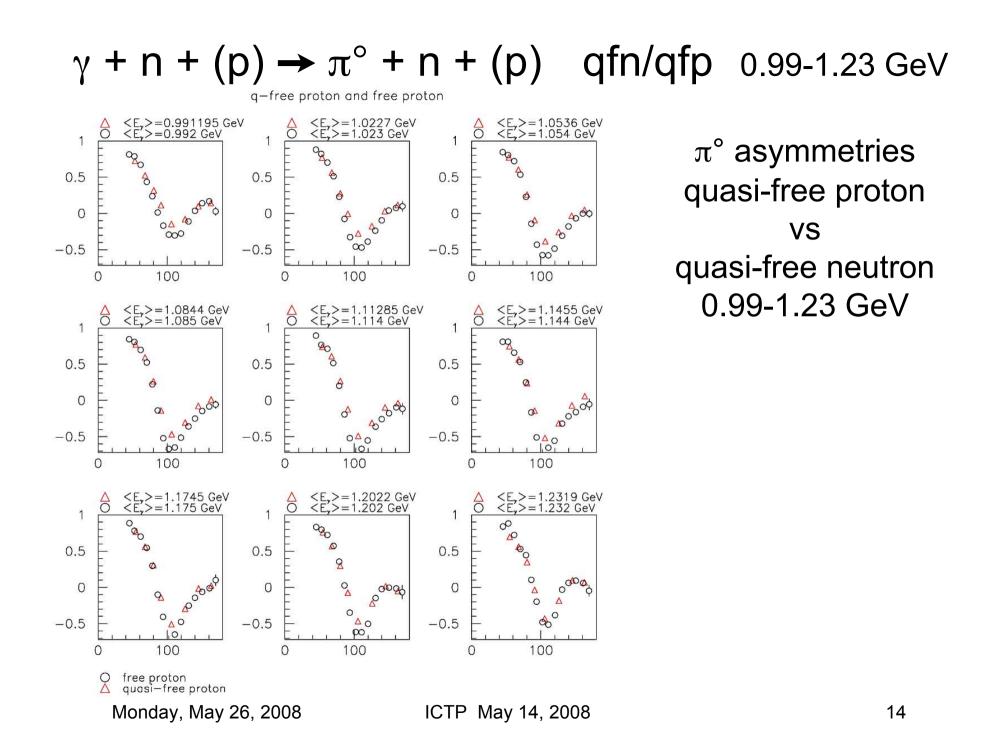


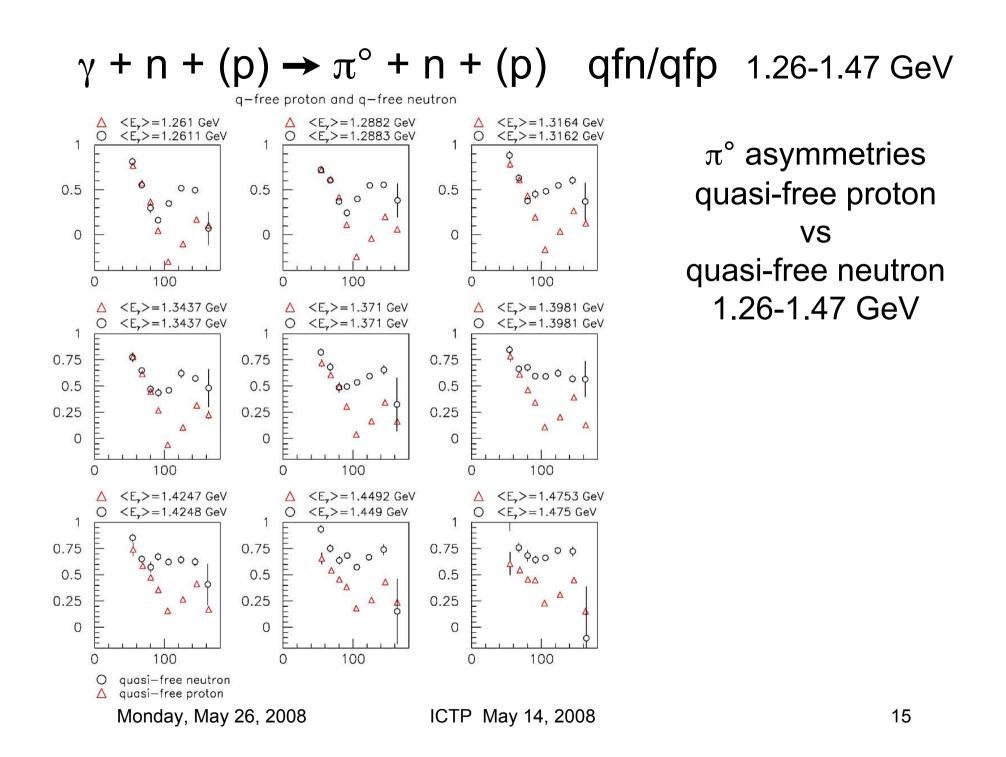


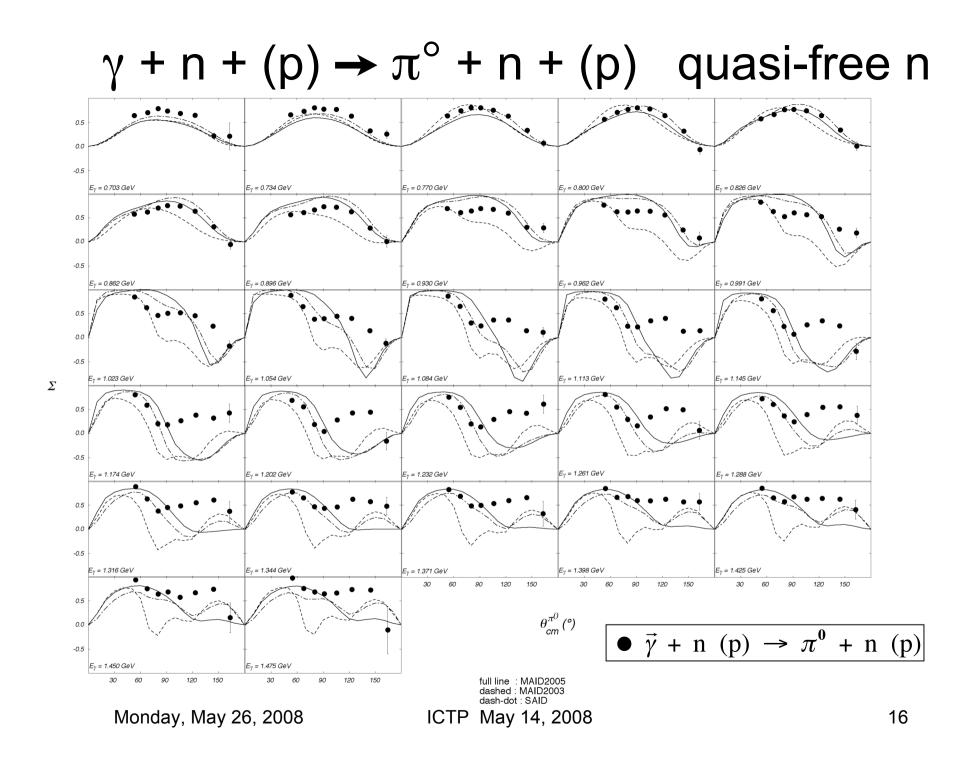












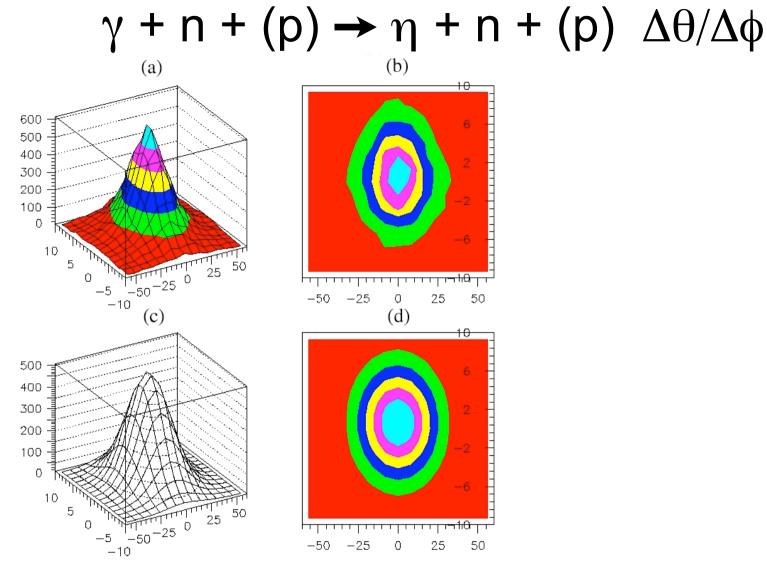


FIG. 1: Upper part: The correlation $\Delta\theta$ vs. $(\Delta\phi-180^{\circ})$ in three-dimensional view(a) and its projection on two dimensions (b); Lower part: The bidimensional gaussian fit (c) and its projection on two dimensions (d).

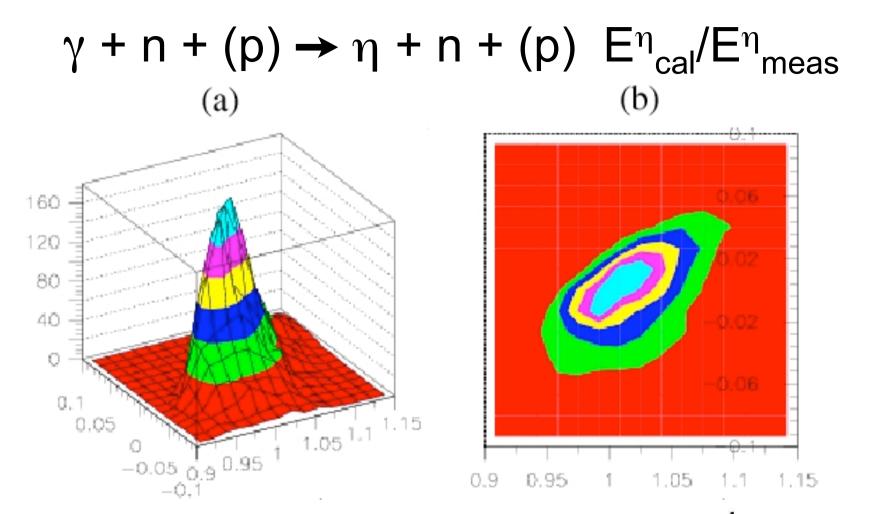
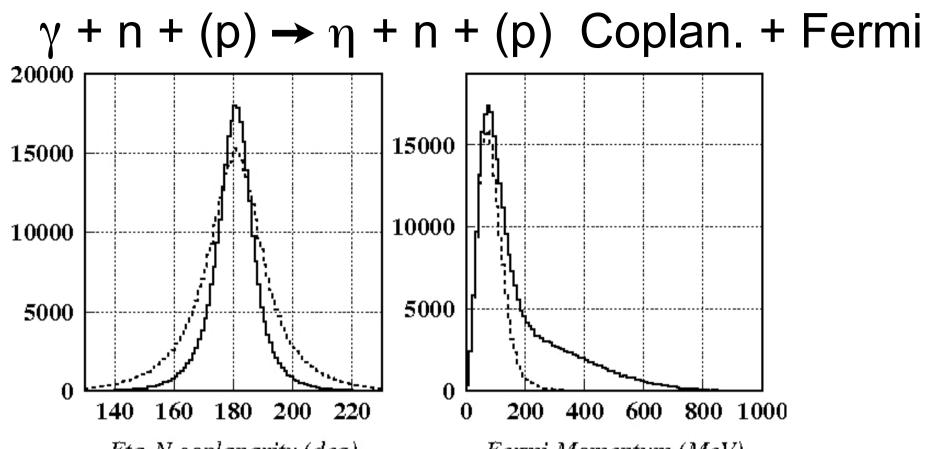


FIG. 2: The three-dimensional correlation $E_{\eta}^{calc}/E_{\eta}^{meas}$ vs. $(M_X - M_N)$ (a) (see text for explanation) and its projection on two dimensions (b).



Eta-N coplanarity (deg) Fermi Momentum (MeV) FIG. 5: Left: Coplanarity between the η and the free proton (solid line) and the quasi-free proton (dotted line): the smearing between the two distributions is due to the Fermi motion; Right: The Fermi momentum distribution before (solid line) and after (dashed line) the application of the bidimensional cuts (see text for details).

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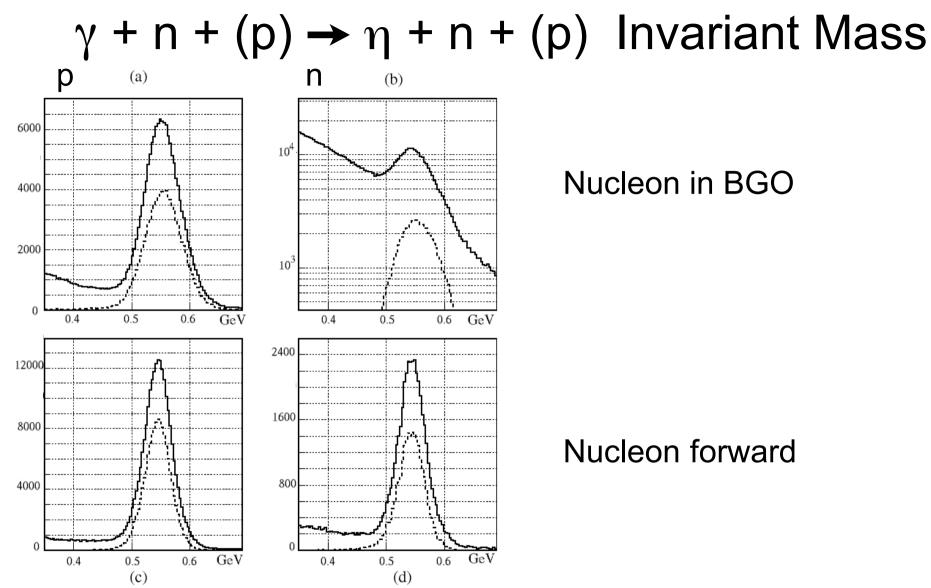
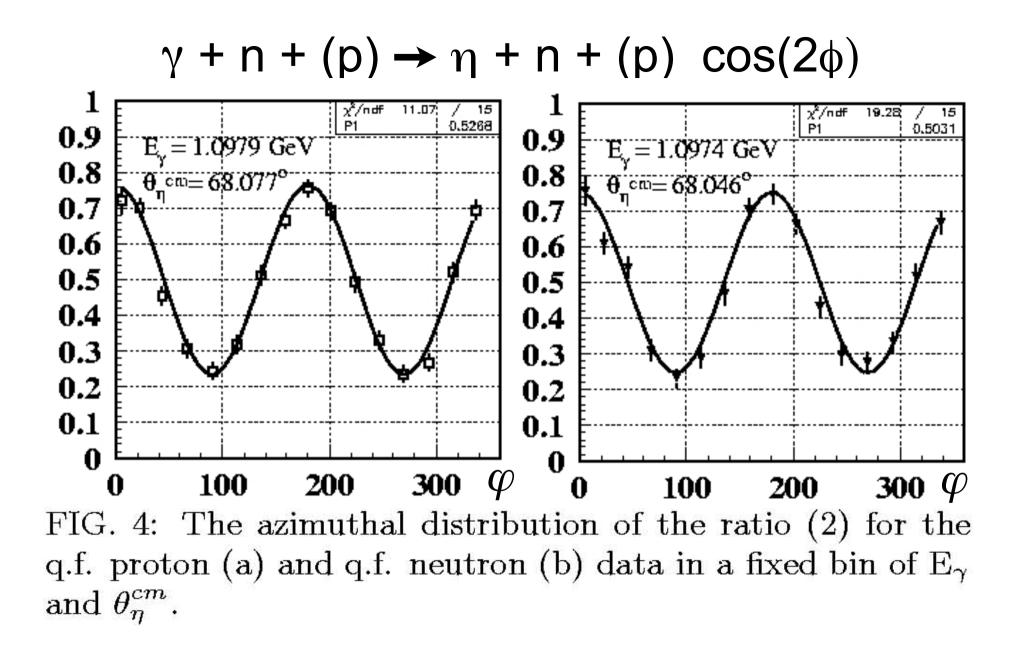


FIG. 3: The η invariant mass without cuts (solid line) and with the kinematical cuts (dotted line) for a central proton (a) and neutron (b) (in logarithmic scale), for a forward proton (c) and neutron (d).



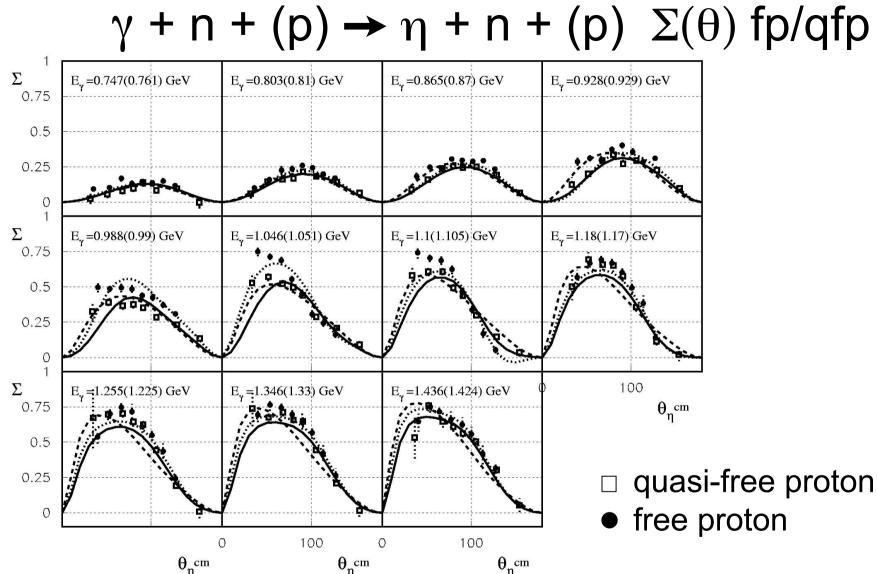


FIG. 6: Beam asymmetry Σ in η photoproduction on the quasi-free proton (open squares) in the deuteron and on the free proton (full circles)[2]. The energy value outside and inside parenthesis indicate the mean value of the bin for quasi-free and free protons respectively. In dotted lines are illustrated the predictions of Maid2001 [7] for the free proton, in solid and dashed lines those for the quasi-free proton of Maid2001 [7] and the reggeized model [3], respectively (see text for details).

$\gamma + n + (p) \rightarrow \eta + n + (p)$ Tiator

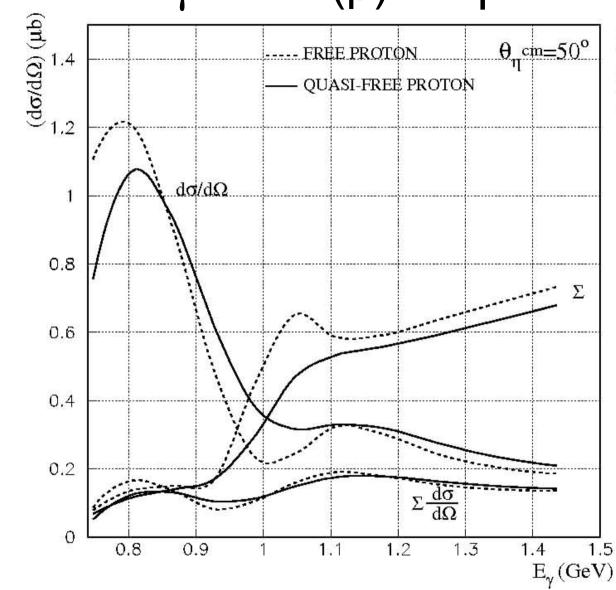


FIG. 7: Energy dependence of the differential cross section for $\gamma p \rightarrow \eta p$ calculated at $\theta = 50^{\circ}$. The unpolarized cross section (5), the polarized cross section (6) and the beam asymmetry Σ are presented in the three pairs of curves for the free proton (dashed curves) and the quasi free proton (solid curves) respectively.

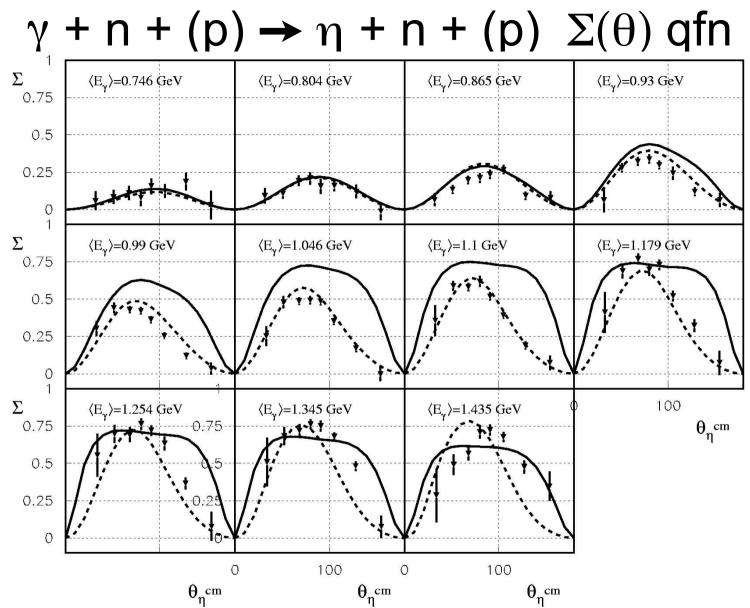


FIG. 8: Beam asymmetry Σ in η photoproduction on the quasi-free neutron in eleven energy bins, plotted as a function of the θ_{η}^{cm} . In each plot, the mean γ energy of the bin is also indicated. In solid and dashed lines are illustrated the predictions for neutrons of Maid2001 and of the reggeized model respectively (see text for details).

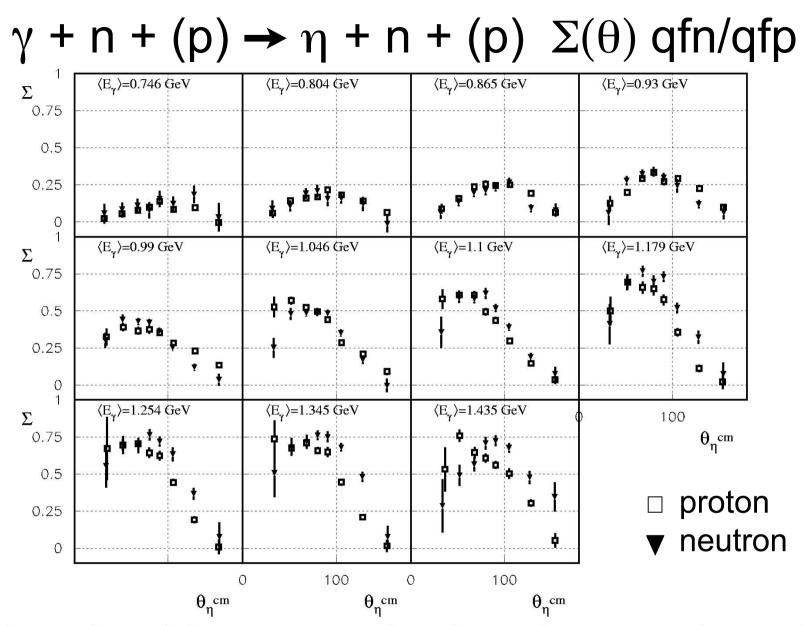
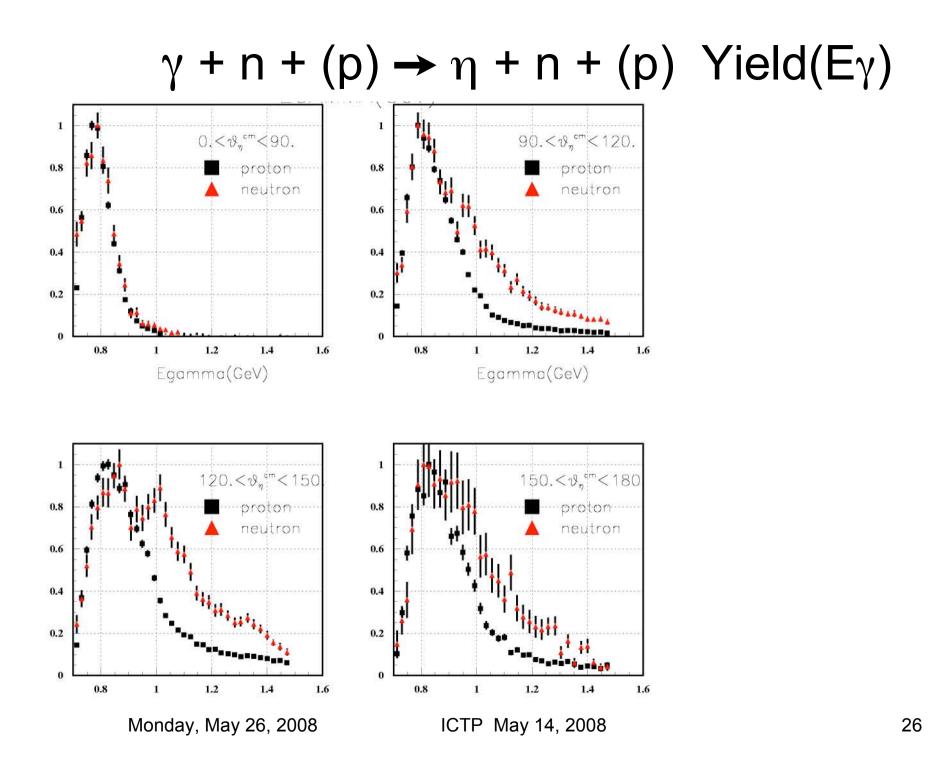
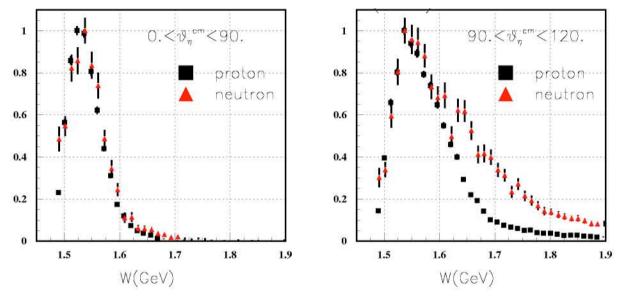
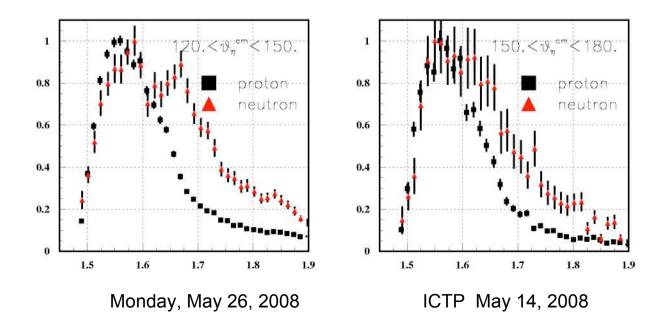


FIG. 9: Comparison between the beam asymmetry Σ in η photoproduction on the quasi-free proton (open squares) and the quasi-free neutron (full triangles) in the eleven energy bins, plotted as a function of the θ_{η}^{cm} . See text for details.



 $\gamma + n + (p) \rightarrow \eta + n + (p)$ Yield(W)





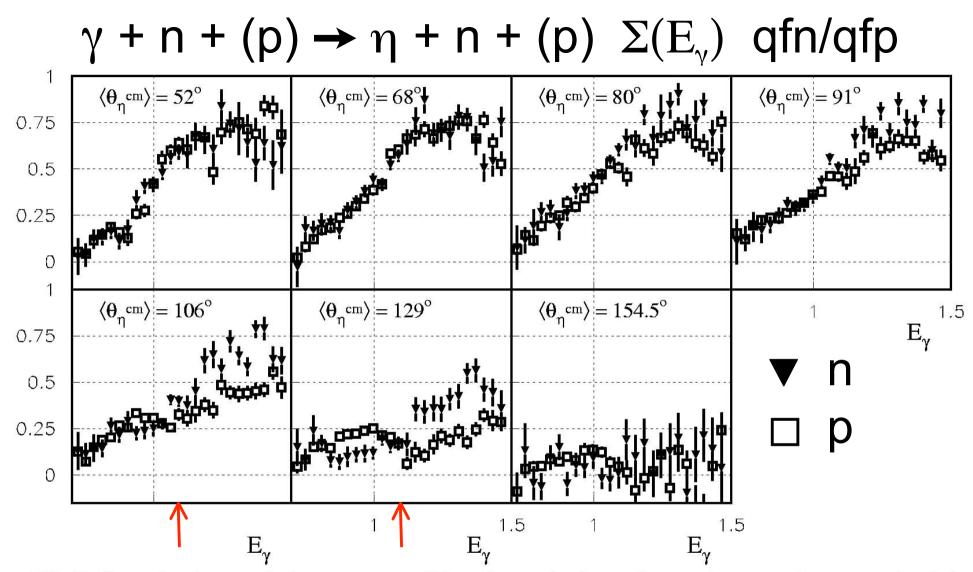


FIG. 10: Comparison between the beam asymmetry Σ in η photoproduction on the quasi-free proton (open squares) and the quasi-free neutron (full triangles) in seven angular bins, plotted as a function of the γ energy (intervals of $\simeq 25$ MeV width.

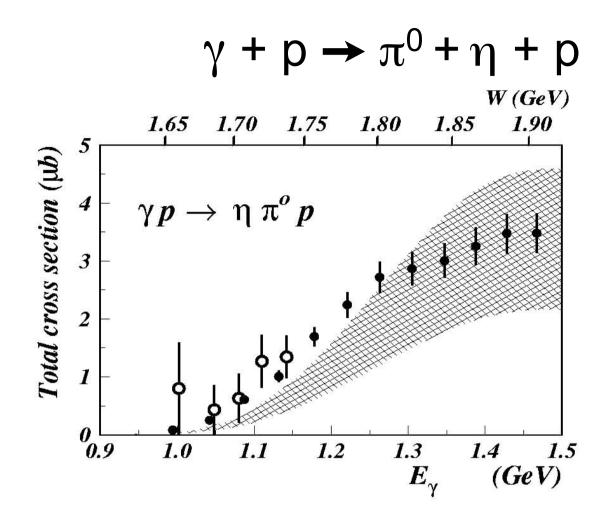
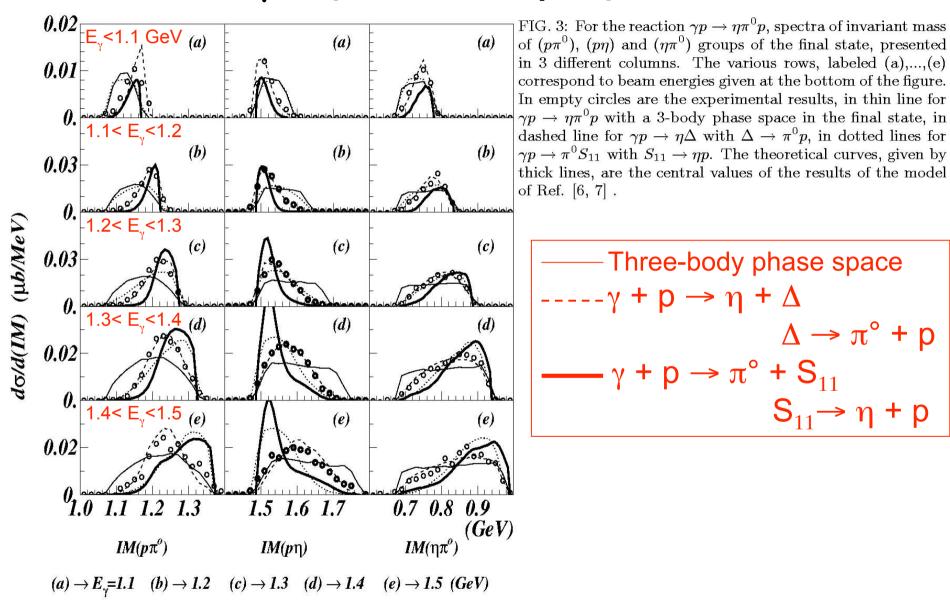


FIG. 2: Total cross section of the reaction $\gamma p \rightarrow \eta \pi^0 p$. The dots are the experimental data of this work. The open circles are from reference [9]. The results of the model of Ref. [6, 7] are given with their uncertainty by a hatched band of the figure. The uncertainty originates from the one on the $\gamma p \Delta(1700)$ coupling which was taken from the PDG [10] Monday, May 26, 2008 ICTP May 14, 2008

 \mathbf{O}_{tot}

 $\gamma + p \rightarrow \pi^0 + \eta + p$ IM



 $\gamma + p \rightarrow \pi^0 + \eta + p \Sigma$

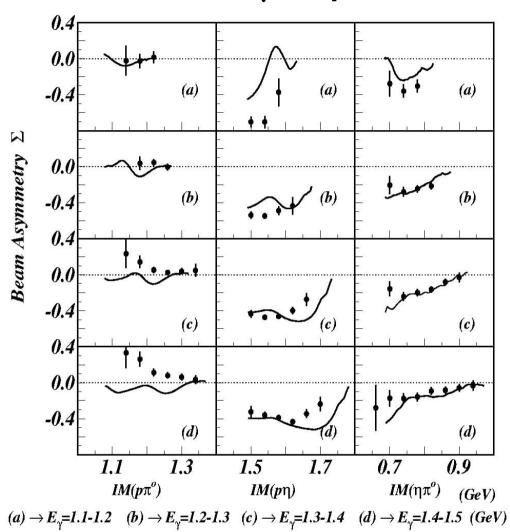


FIG. 4: Beam asymmetry of the reaction $\gamma p \rightarrow \eta \pi^0 p$. The theoretical results are calculated with the model of Ref. [6, 7]

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K- Λ e K- Σ

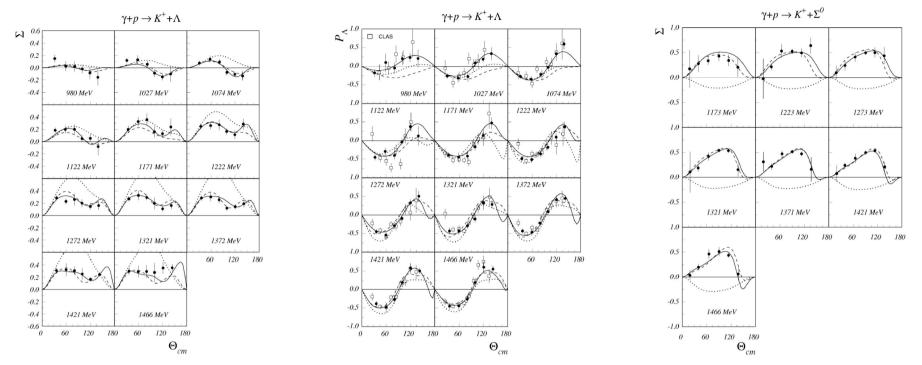


Fig. 11. Angular distributions of the beam asymmetries Σ for $\gamma p \rightarrow K^+ A$ and γ -ray energies ranging from threshold up to 1500 MeV. Data are compared with the Bonz0005 solutions with large (dashed lines) and small (solid lines) S₁₁ contributions 1500 MeV. GRAAL data (closed circles) are compared with CLAS data (open squares). Definition of the curves as in fig. 11. 1500 MeV. Definition of the curves as in fig. 11.

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Graal closed circles

CLAS open squares

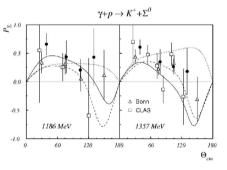
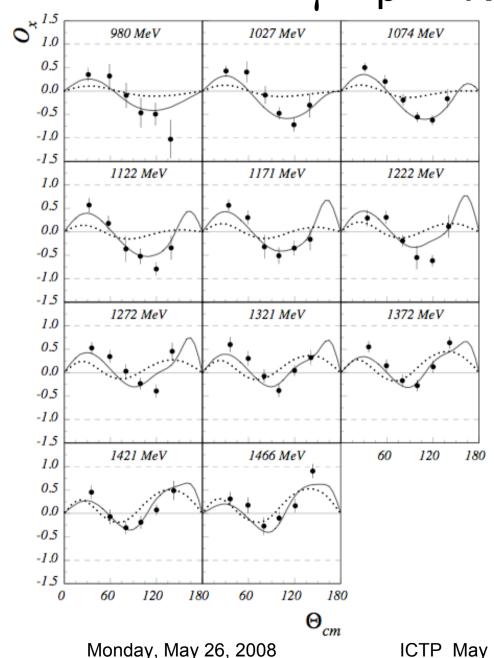
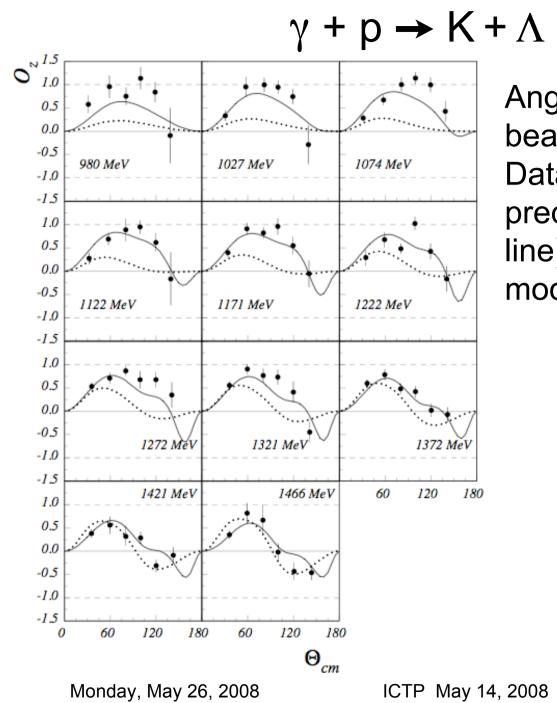


Fig. 10. Σ^0 recoil polarizations for $\gamma p \to K^+ \Sigma^0$. Comparison between GRAAL (closed circles), Bonn (open triangles) and CLAS (open squares) data. Data are compared with the Bonn2005 solutions with large (dashed lines) and small (solid lines) S₁₁ contribution at low energies and with the Kaon-MAID 2000 standard calculation (dotted lines).

 $\gamma + p \rightarrow K + \Lambda$ $O_{\mathbf{x}}$

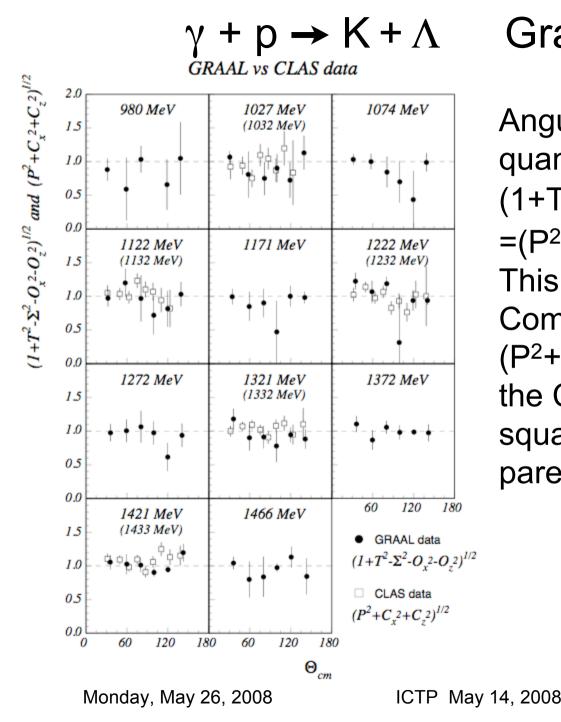


Angular distributions of the beam recoil observable O_x . Data are compared with the predictions of two models: solid line BCC (Bonn Coupled Channel - A. V. Anisovich et al. Eur. Phys. J. A 25, 427 (2005), A. V. Sarantsev et al. Eur. Phys. J. A 25, 441(2005)); dotted line GRPR (Ghent Regge Plus Resonance - T. Corthals, J. Ryckerbush and T. Van Cauteren, Phys. Rev. C 73, 045207 (2006)).



Angular distributions of the beam recoil observable O_z . Data are compared with the predictions of the BCC (solid line) and GRPR (dotted line) models.

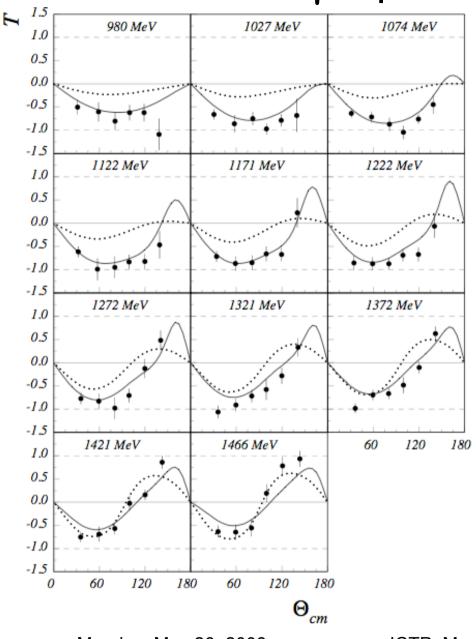
 O_7



Graal vs CLAS

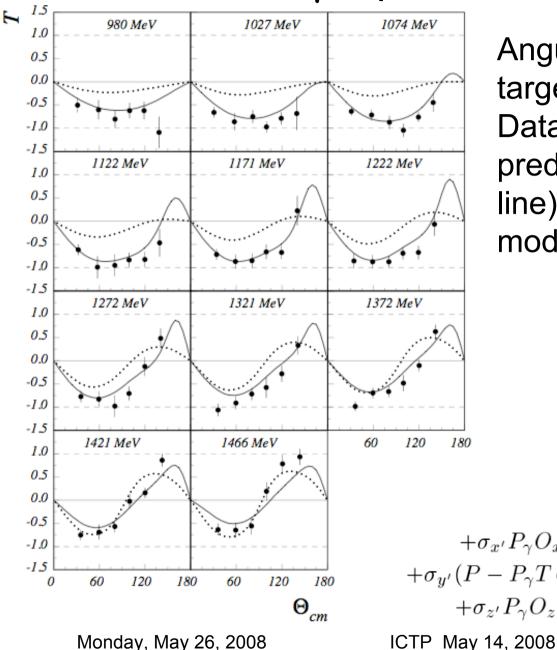
Angular distributions of the quantity $(1+T^2-\Sigma^2-O_x^2-O_z^2)^{1/2} =$ $=(P^2+C_x^2+C_z^2)^{1/2}$. This quantity should be ≤ 1 . Comparison to the values $(P^2+C_x^2+C_z^2)^{1/2}$ published by the CLAS collaboration (open squares - energy in parentheses).





Angular distributions of the target asymmetry T. Data are compared with the predictions of the BCC (solid line) and GRPR (dotted line) models.





Angular distributions of the target asymmetry T. Data are compared with the predictions of the BCC (solid line) and GRPR (dotted line) models.

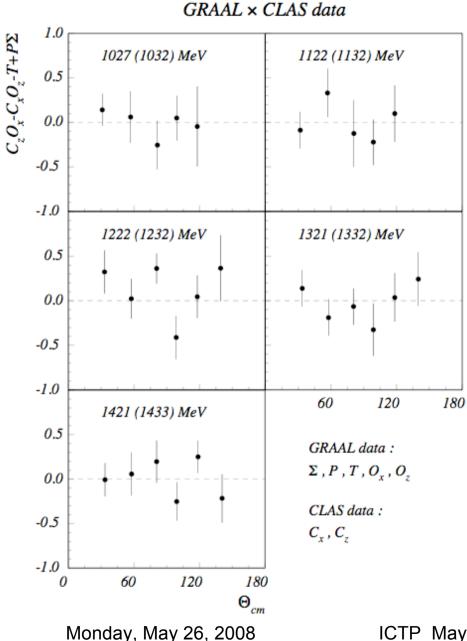
$$\rho_f \frac{d\sigma}{d\Omega} = \frac{1}{2} \left(\frac{d\sigma}{d\Omega} \right)_0 [1 - P_\gamma \Sigma \cos 2\varphi_\gamma]$$

$$= \frac{1}{2} \left(\frac{d\sigma}{d\Omega} \right)_0 [1 - P_\gamma \Sigma \cos 2\varphi_\gamma]$$

$$= -P_\gamma T \cos 2\varphi_\gamma]$$

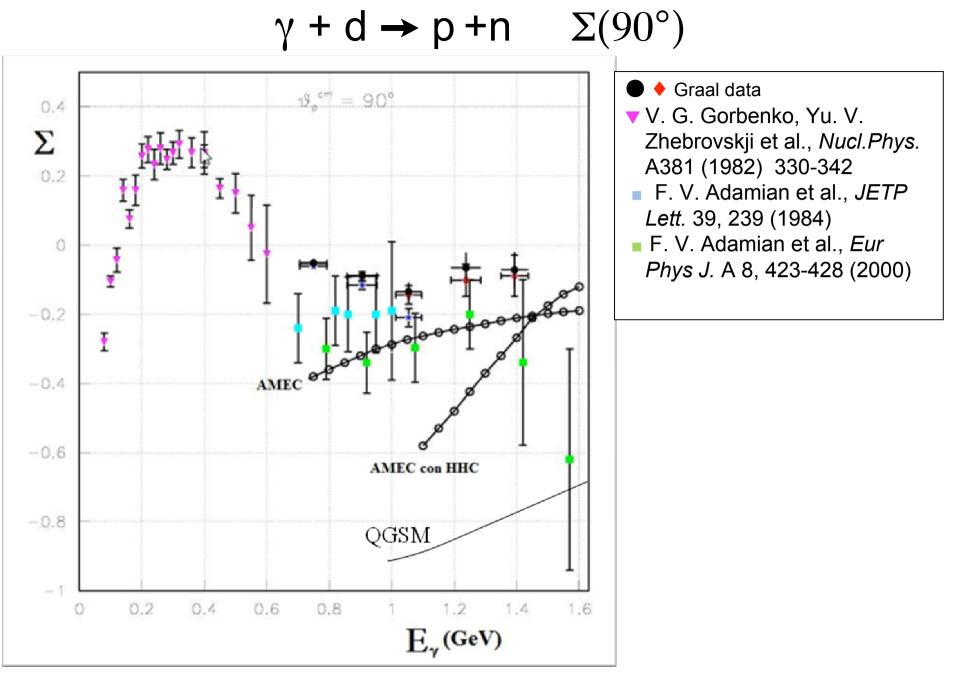
$$= \frac{1}{2} \left(\frac{1}{2} \left(\frac{1}{2} - \frac{1}{2} \right)_0 \left(\frac{1}{2} - \frac{1}{2} \right)_0 \left(\frac{1}{2} - \frac{1}{2} \right)_0 \left(\frac{1}{2} - \frac{1}{2} - \frac{1}{2} - \frac{1}{2} \right)_0 \left(\frac{1}{2} - \frac{1}{2} - \frac{1}{2} - \frac{1}{2} \right)_0 \left(\frac{1}{2} - \frac{1}{2} - \frac{1}{2} - \frac{1}{2} \right)_0 \left(\frac{1}{2} - \frac{1}{2}$$

$\gamma + p \rightarrow K + \Lambda$



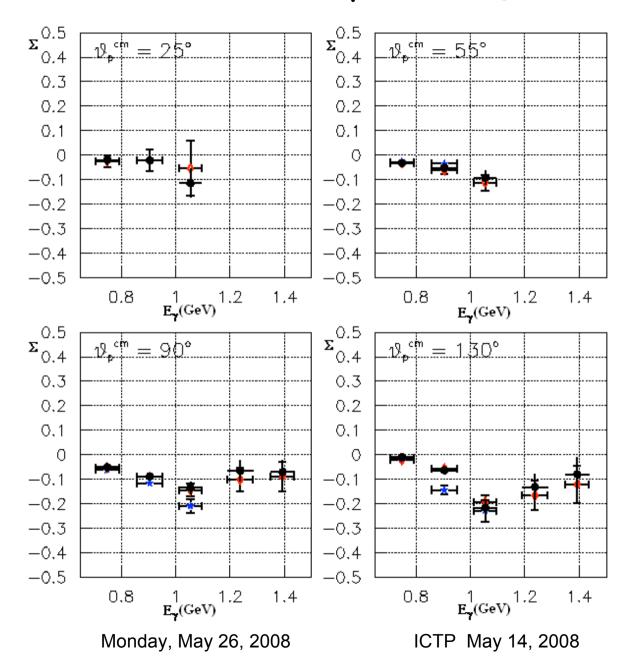
Graal and CLAS

Angular distributions of the quantity $C_7 O_7 - C_7 O_7 - T + P \Sigma$. It is calculated using the C_x and C_z results published by the CLAS collaboration (energy in parentheses) combined with our O_x and O_z data converted to have the same z' axis convention and with our Σ , P and T measurements. The used CLAS data are those corresponding to the angles $\cos(\theta_{cm})=0.85$, mean(0.65,0.45), mean(0.25,0.05), -0.15, mean(-0.35,-0.55) and -0.75. We should have the equality $C_{z} O_{x} - C_{y} O_{z} - T + P \Sigma = 0$



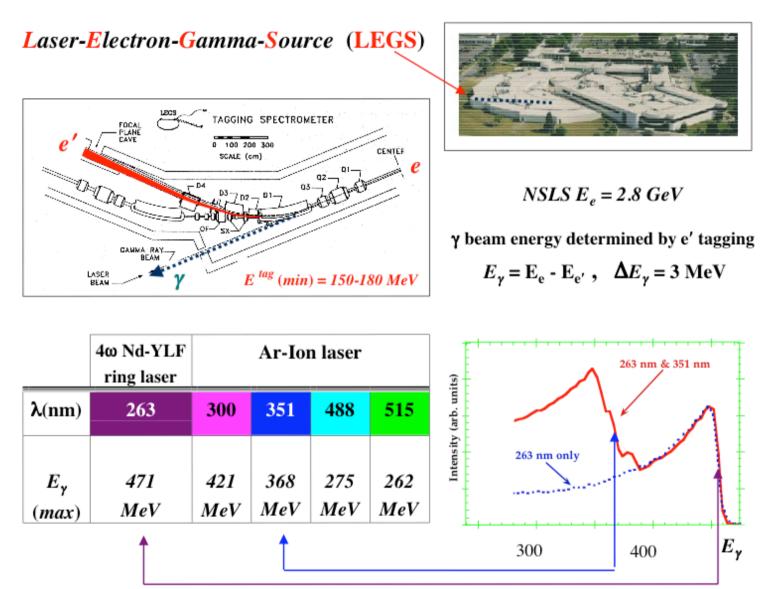
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Σ $\gamma + d \rightarrow p + n$



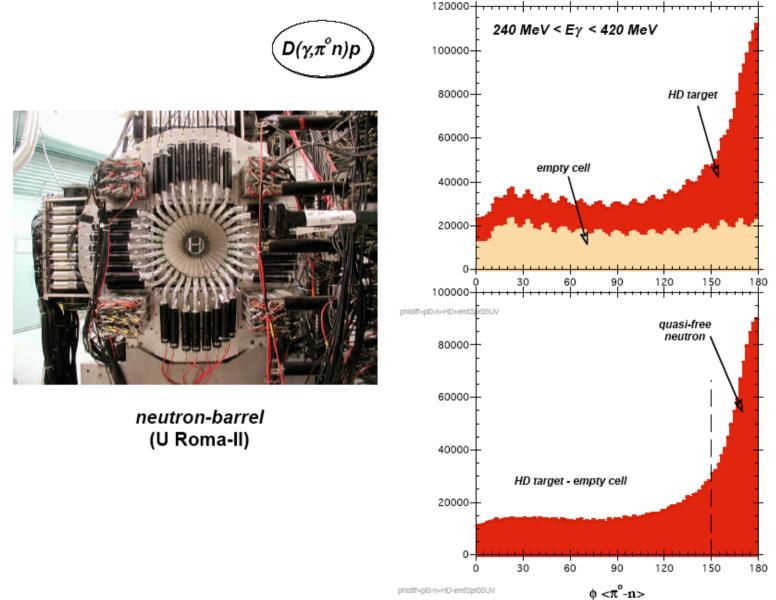
Graal data analyzed in different ways

LEGS at BNL - Summary



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LEGS at BNL: Spring 05 Neutron Detection



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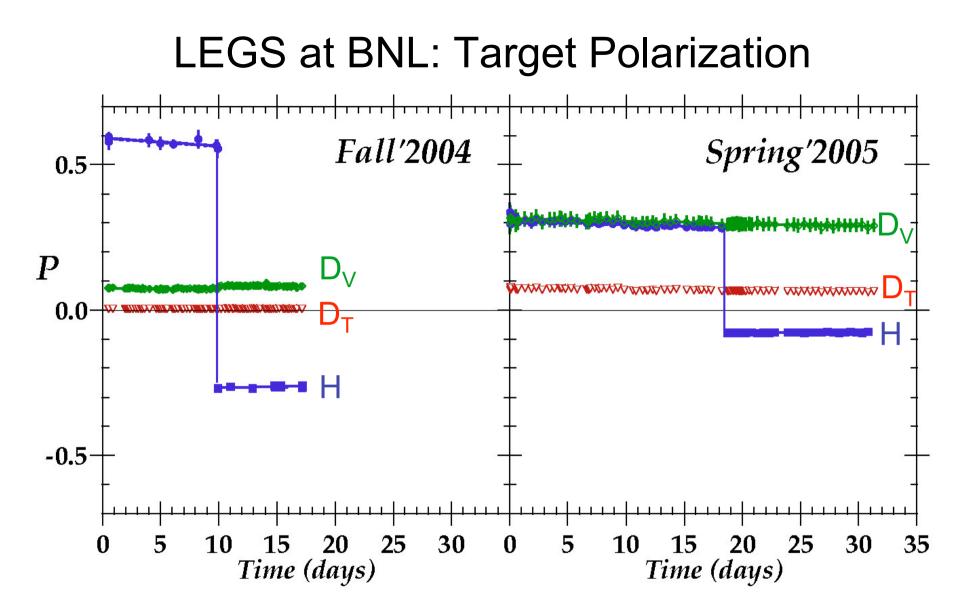
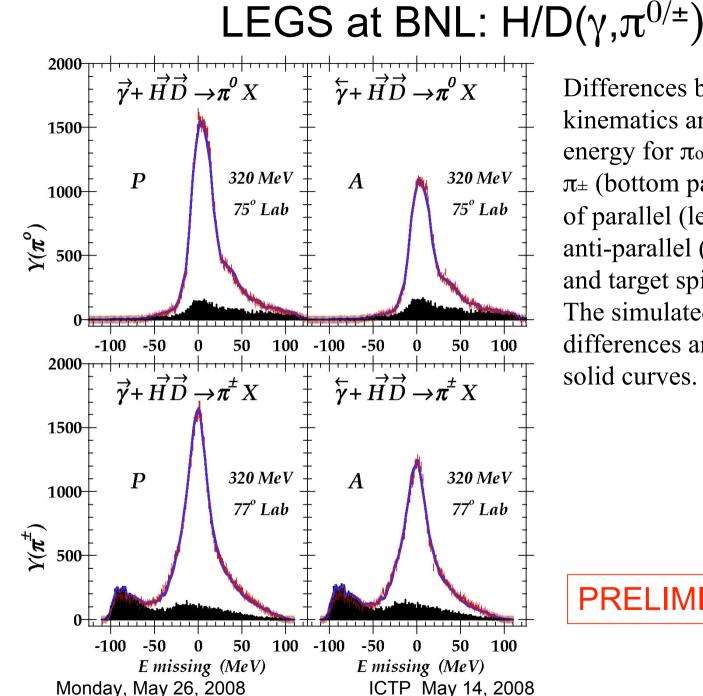
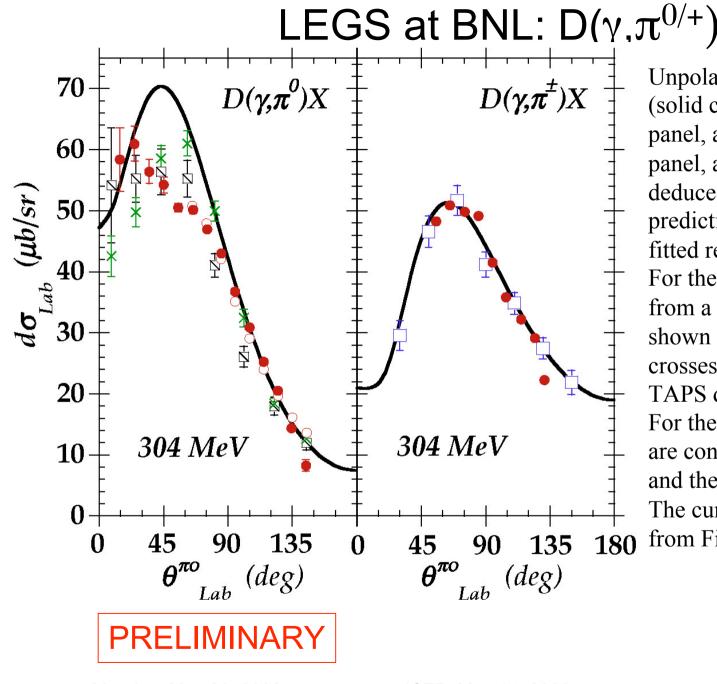


Figure 1. Polarizations of H (blue) and D (green-vector and red-tensor) nuclei in HD during the two data collection periods. Mid-way through each, the H polarization was flipped using an RF transition.



Differences between 2-body kinematics and the measured energy for π_0 (top panels) and π [±] (bottom panels), in the cases of parallel (left panels) and anti-parallel (right panels) beam and target spin alignments. The simulated energy differences are shown as the solid curves.





Unpolarized cross sections (solid circles) for $D(\gamma,\pi_0)X$, left panel, and D(γ , π [±])X, right panel, at $E_{\gamma} = 304$ MeV, deduced by subtracting SAID predictions[8] for $p(\gamma, \pi)$ from the fitted results for HD. For the π_0 channel, LEGS data from a liquid D₂ target are shown as open circles, while crosses and hatched-boxes are TAPS data from [18] and [17]. For the π_{\pm} channel, open boxes are constructed from π -pp [19] and the π -/ π + ratio data of [20]. The curves are calculations from Fix and Arenhövel [12].

LEGS at BNL: $H(\gamma, \pi^{0/+})$

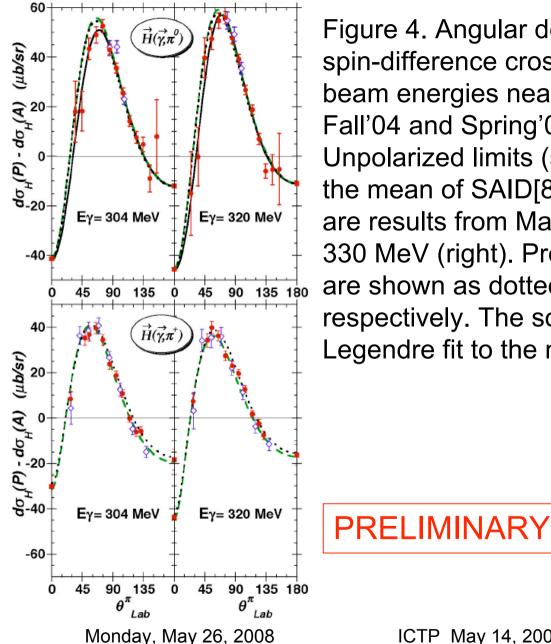


Figure 4. Angular dependence of the $[d\sigma(P) - d\sigma(A)]$ spin-difference cross section for polarized H at beam energies near the Δ peak. The full data from Fall'04 and Spring'05 are shown as solid circles. Unpolarized limits (solid squares) at 0o and 180o are the mean of SAID[8] and MAID[9]. Open diamonds are results from Mainz [21] at 310 MeV (left) and at 330 MeV (right). Predictions from SAID and MAID are shown as dotted and dashed curves, respectively. The solid curves in the top panels are a Legendre fit to the new data.

Compton Scattering Kinematics

2. The maximum energy lost by the electrons after an elastic scattering with a laser photon is given by the maximum energy acquired by the photon:

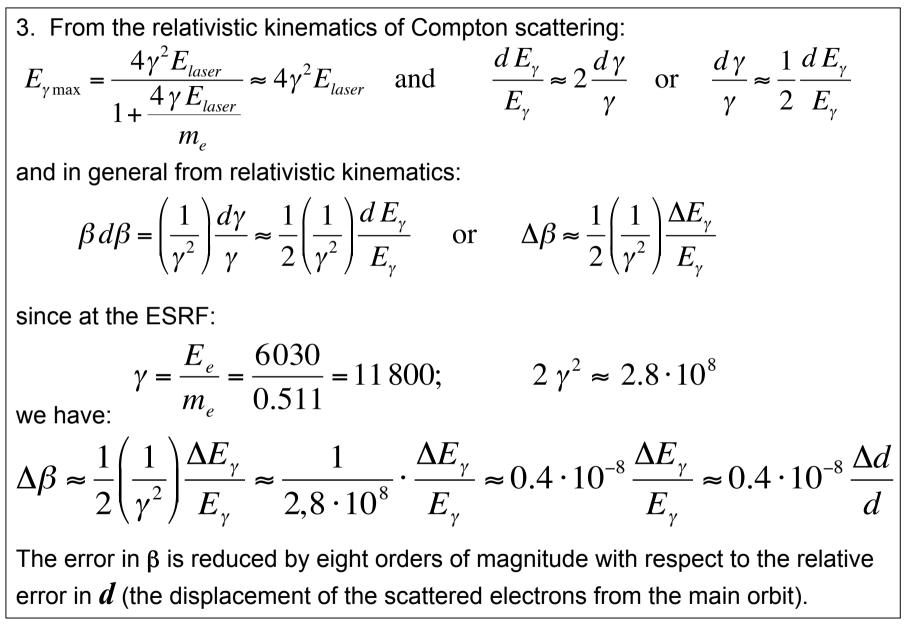
$$E_{el}^{0} - E_{el}^{scatt} = E_{\gamma \max} = \frac{4\gamma^{2}E_{laser}}{1 + \frac{4\gamma E_{laser}}{m_{e}}} \approx 4\gamma^{2}E_{laser}$$

This energy loss is measured by the displacement **d** of the scattered electrons from the primary electron beam after the first magnetic dipole. For the ESRF electron energy of 6.03 GeV and a UV laser line of 3.53 eV, the energy loss is 1.487 GeV and corresponds to an electron displacement at the position of the Graal tagging detector: $d \approx 52.3$ mm.

The microstrips of the Graal tagging detector measure the displacement d of the scattered electrons from the main orbit and therefore the energy lost by the electrons (and acquired by the gamma-rays):

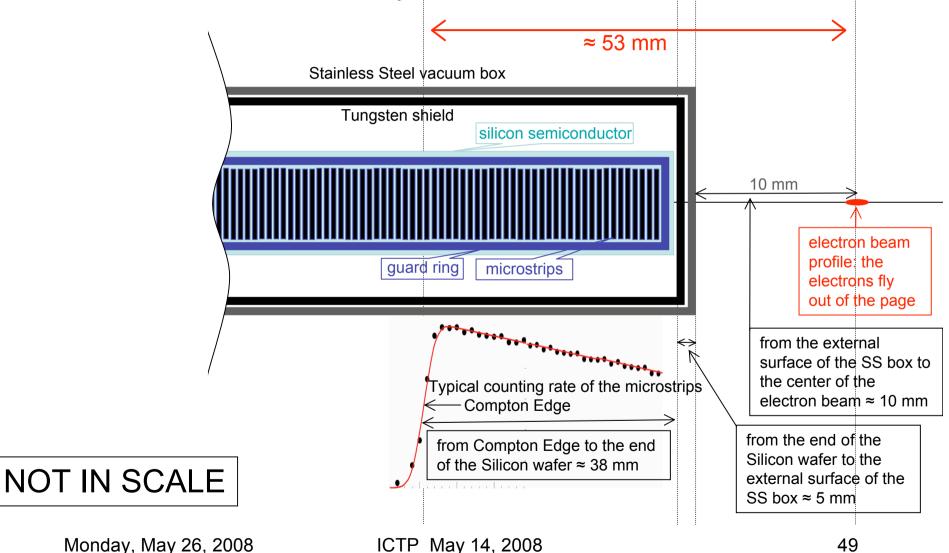
 $E_{\gamma} \propto d$

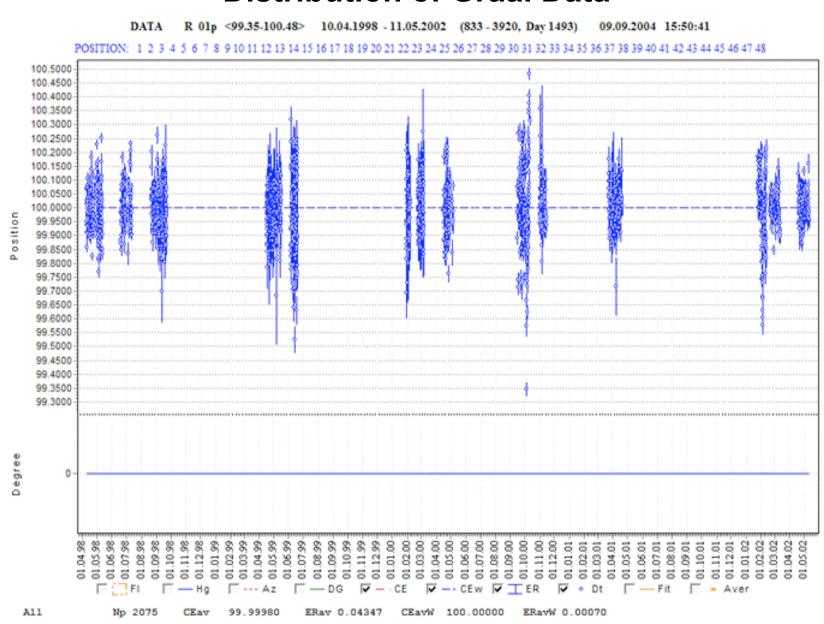
Compton Scattering Kinematics



Graal Tagging Microstrips

Schematic description of the tagging detector in more details. Vertical cut: the electrons fly out of the screen.

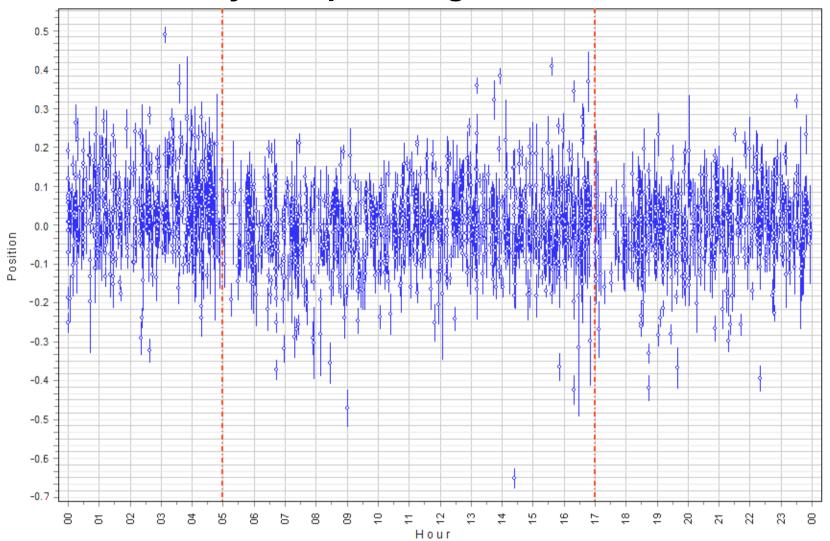




Distribution of Graal Data

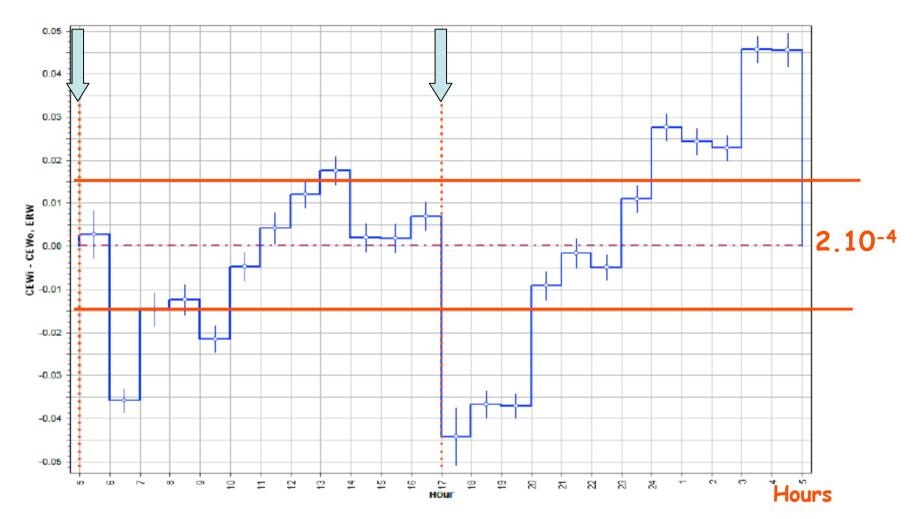
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Daily Compton Edges Distributions



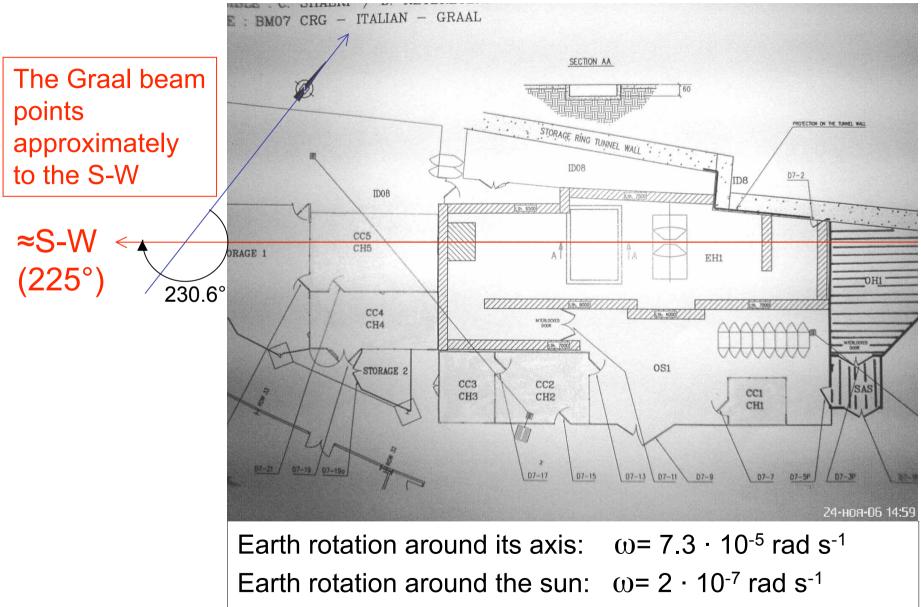
Experimental data plotted as a function of (solar) hour, showing their daily variation. The dispersion of data around the average, taken arbitrarily at zero, is expressed in fractions of microstrip (300 micrometers width or about 7 MeV for one microstrip).

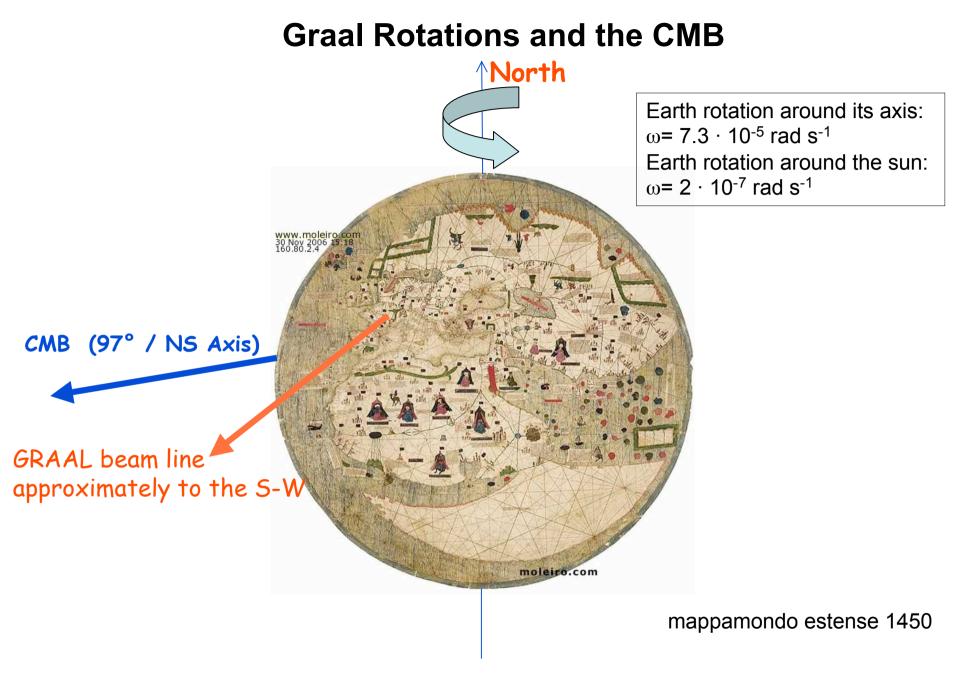
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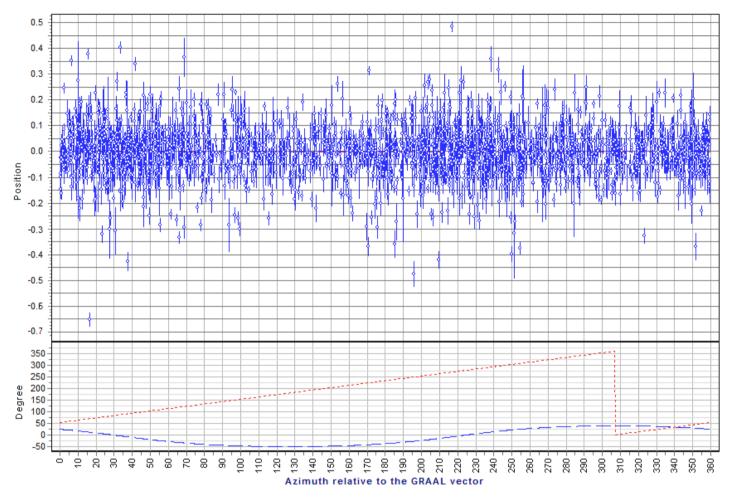
Same as the previous figure, but each point is the average over one hour. The dotted lines show the refill time of the machine corresponding to a possible change in the temperature of the tagging detector or the position of the beam. The average is expressed in microstrip fractions (0.01 = 3μ m).

Graal Beam Orientation on the Earth





Compton Edge Positions vs CMB Dipole



Experimental data plotted as a function of the azimuth (above); below, the variation of the angle between the beam and the CMB dipole decomposed to azimuth (dotted) and declination (dashed) angles is shown.

Assuming an error of

2 · 10⁻⁴

in our determination of the position of the Compton Edge, we could arrive to an estimated upper limit on the asymmetry of the velocity of light of:

$$\Delta\beta \approx \frac{1}{2} \left(\frac{1}{\gamma^2}\right) \frac{\Delta E_{\gamma}}{E_{\gamma}} \approx 0.4 \cdot 10^{-8} \frac{\Delta d}{d} \approx 0.4 \cdot 10^{-8} \cdot 2 \cdot 10^{-4} \approx 10^{-12}$$

Considering that we have analyzed old data and we have not been able to reconstruct completely the status of the system - accelerator + tagging detector - during our runs we have published the more conservative number:

3 · 10⁻¹²

An Optimistic View of the Future

In conclusion if optimistically we assume a systematic error of 2.5 μ m in the distance between the position of the Compton Edge and the electron beam, we have:

$$\frac{\left(\Delta d\right)_{sys}}{d} \approx \frac{2.5\,\mu m}{52.3\,mm} \approx 5\cdot 10^{-5} \approx \frac{\left(\Delta E_{\gamma}\right)_{stat}}{E_{\gamma}}$$

and we can hope to be able to verify the isotropy of the velocity of light with respect to some absolute reference frame with a precision of:

$$\Delta\beta \approx \frac{1}{2} \left(\frac{1}{\gamma^2}\right) \frac{\Delta E_{\gamma}}{E_{\gamma}} \approx 0.4 \cdot 10^{-8} \frac{\Delta E_{\gamma}}{E_{\gamma}} \approx 0.4 \cdot 10^{-8} \cdot 5 \cdot 10^{-5} \approx 2 \cdot 10^{-13}$$