Unquenching the Quark Model.

E. Santopinto and R. Bijker

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Many versions of CQMs have been developed (KI, CI, GBE, U(7), hCQM,Bonn, etc.) non relativistic and relativistic

CQMs:

Good description of the spectrum Predictions of many quantities: photocouplings helicity amplitudes elastic form factors structure functions Based on the effective degrees of freedom of 3 constituent quarks There are Open Problems

- Missing resonances
- Low Q² behavior of helicity amplitudes
- etc. ...

Is it a degrees of freedom problem?

$q\overline{q}$ corrections ?





QM's problems ----> degrees of freedom

Key degrees of freedom---->qq pair



Problems

1) find a quark pair creation mechanism QCD inspired

2) implementation of this mechanism in such a way to

do not destroy the good QMs results

Isgur's flux-tube-breaking model for mesons

In the flux tube model Isgur has shown that a "miraculous" set of cancellations between apparently uncorrelated sets of mesons occurs in such a way that

-the OZI hierarchy is preserved

-there is near immunity of the long range potential

it is necessary a to sum over very large towers of intermediate states to see that the spectrum is only weekly perturbed great shift only from non adiabatic effects, as the result of the coupling of a resonance to a very nearby threshold for decays into other two.

quark potential model as adiabatic limit

Unquenched Quark Model



Strange quark-antiquark pairs in the proton with h.o. wave functions

Geiger & Isgur, PRD 55, 299 (1997)

- Pair-creation operator with ³P₀ quantum numbers of vacuum and geometry of flux-tube breaking model
- Sum over a large tower of intermediate states to preserve the phenomenological success of CQM's

The creation mechanism is at the quark level -> (b)



 Σ over all the towers of intermediate states

Pair creation model applied by Isgur to the strangeness content of the proton, or $O_s = \Delta S$, R_s^2, μ_s

$$|p\rangle \to |p\rangle + \sum_{Y^*K^*\ell S} \int q^2 dq \ |Y^*K^*q\ell S\rangle \frac{\langle Y^*K^*q\ell S |h_{s\bar{s}}|p\rangle}{M_p - E_{Y^*} - E_{K^*}} ,$$

$$\begin{array}{lcl} \langle O_s \rangle & = & \sum_{\substack{Y^* K^* \ell S \\ Y^{*'} K^{*'} \ell' S'}} \int q^2 dq \; q'^2 dq' \; \frac{\left\langle p \; | h_{s\bar{s}} | \; Y^{*'} K^{*'} q' \ell' S' \right\rangle}{M_p - E_{Y^{*'}} - E_{K^{*'}}} \\ & \quad \times \left\langle Y^{*'} K^{*'} q' \ell' S' \; | O_s | \; Y^* K^* q \ell S \right\rangle \frac{\left\langle Y^* K^* q \ell S \; | h_{s\bar{s}} | \; p \right\rangle}{M_p - E_{Y^*} - E_{K^*}} \end{array}$$



$$\begin{array}{lll} \langle O_s \rangle &=& \sum_{\substack{Y^* K^* \ell S \\ Y^{*'} K^{*'} \ell' S'}} \int q^2 dq \; q'^2 dq' \; \frac{\left\langle p \; |h_{s\bar{s}}| \; Y^{*'} K^{*'} q' \ell' S' \right\rangle}{M_p - E_{Y^{*'}} - E_{K^{*'}}} \\ &\times \left\langle Y^{*'} K^{*'} q' \ell' S' \; |O_s| \; Y^* K^* q \ell S \right\rangle \frac{\left\langle Y^* K^* q \ell S \; |h_{s\bar{s}}| \; p \right\rangle}{M_p - E_{Y^*} - E_{K^*}} \end{array}$$

Closure limit ---->

$$\langle O_s \rangle \propto \langle p \left| h_{s\bar{s}} O_s h_{s\bar{s}} \right| p \rangle \propto \langle 0 \left| h_{s\bar{s}} O_s h_{s\bar{s}} \right| 0 \rangle \ ,$$

Good stringent test for the program. The corrections due to qq-pairs are zero in the closure limit. The sums over towers of states are constrained by the closure limit in such a way that very different meson and baryon states compensate.

It would be desirable to devise tests of the mechanisms underlying the delicate cancellations which conspire to hide the effects of the sea in the picture presented here. It also seems very worthwhile to extend this calculation to uu and dd loops. Such an extension could reveal the origin of the observed violations [38] of the Gottfried sum rule [39] and also complete our understanding of the origin of the spin crisis. From our previous calculations [4], the effects of "un-

Geiger & Isgur, PRD 55, 299 (1997)

Extensions

- To any initial baryon or baryon resonance
- To any flavor of the quark-antiquark pair
- To any model of baryons and mesons

Santopinto & Bijker, nucl-th/0701227 Bijker & Santopinto, nucl-th/0703053

Problems for the baryons---->

 big towers of states authomatically generated by means of group theoretical methods

problems linked with permutational symmetry(many different diagrams)-> solved with group theoretical methods

$$|\psi_A\rangle = \mathcal{N} \left\{ |A\rangle + \sum_{qBClJ} \int d\vec{k} |BC\vec{k}lJ\rangle \frac{\langle BC\vec{k}lJ | h_{q\bar{q}}^{\dagger} |A\rangle}{M_A - E_B - E_C} \right\}$$

$$\mathcal{O} = \langle \psi_A \mid \hat{\mathcal{O}} \mid \psi_A \rangle = \mathcal{O}_{\text{valence}} + \mathcal{O}_{\text{sea}}$$

$$\mathcal{O}_{\text{valence}} = \mathcal{N}^2 \langle A \mid \hat{\mathcal{O}} \mid A \rangle$$

$$\mathcal{O}_{\text{sea}} = \mathcal{N}^2 \sum_{qBClJ} \int d\vec{k} \sum_{q'B'C'l'J'} \int d\vec{k}' \frac{\langle A \mid h_{q'\bar{q}'} \mid B'C'\vec{k}'l'J' \rangle}{M_A - E_{B'} - E_{C'}}$$

$$\langle B'C'\vec{k}'l'J' \mid \mathcal{O} \mid BC\vec{k}lJ \rangle \frac{\langle BC\vec{k}lJ \mid h_{q\bar{q}}^{\dagger} \mid A \rangle}{M_A - E_B - E_C}$$

The closure-limit and the symmetries give strong constraints





In the closure limit $\Delta u: \Delta d: \Delta s = 4:-1:0$

Conclusions

- Unquenched quark model:we have extended it to any flavor
- First results very promising
- Future: applications to many observables.